

On the Cause of Solar Differential Rotation

太陽差動自轉的成因

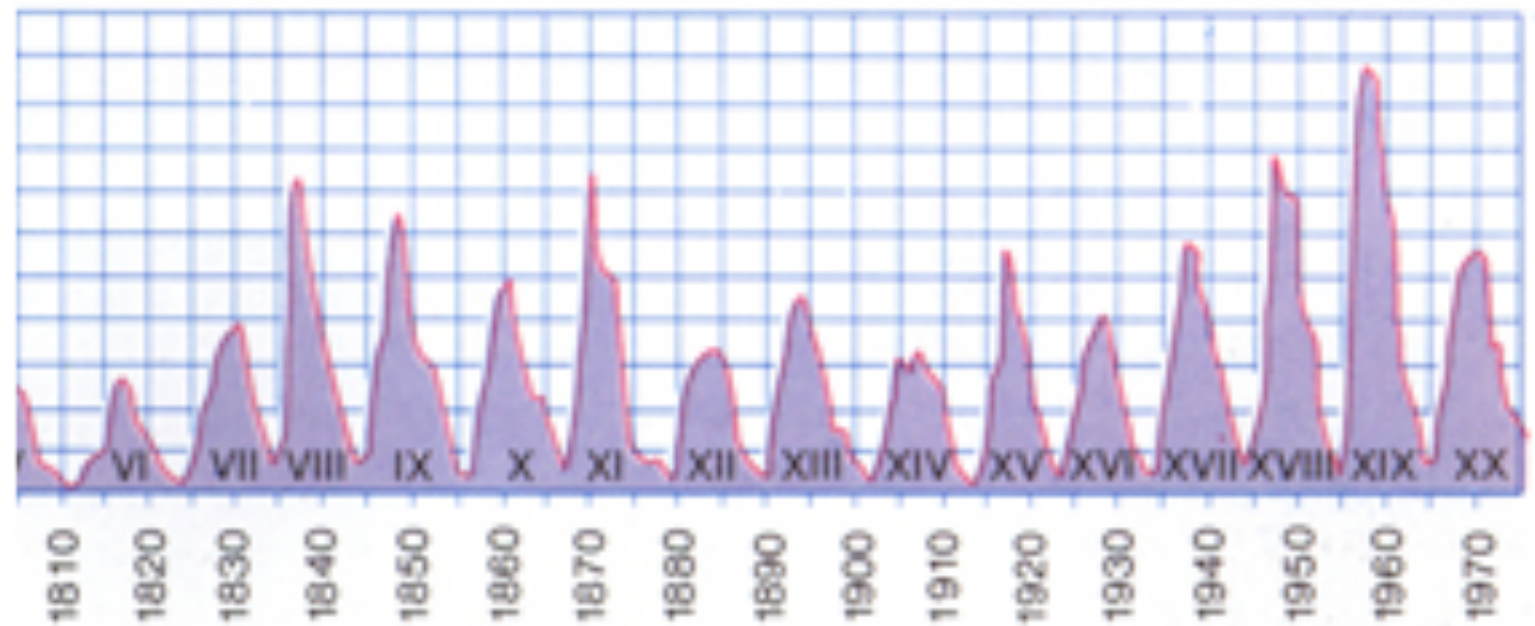
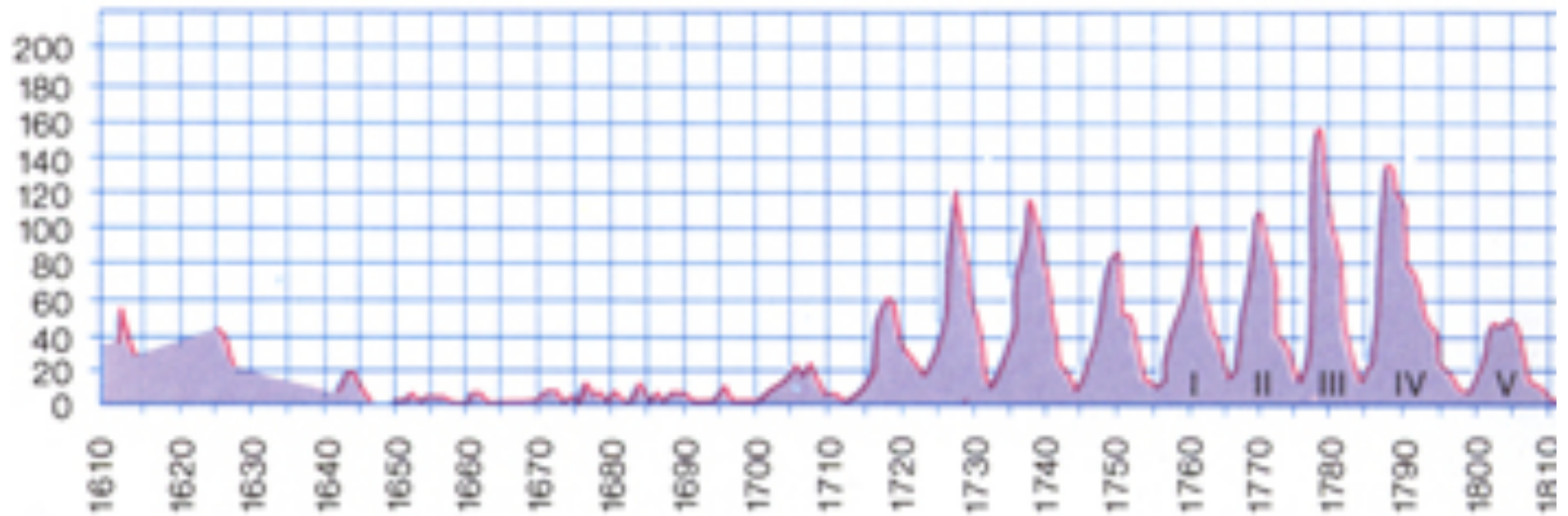
Ling-Hsiao Lyu

Institute of Space Science, National Central University

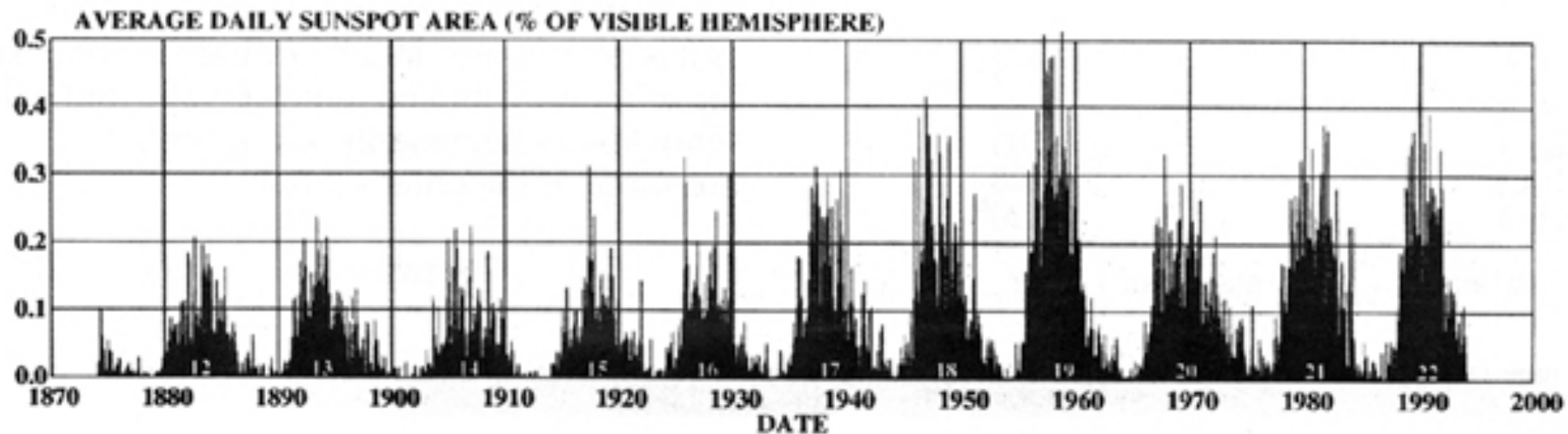
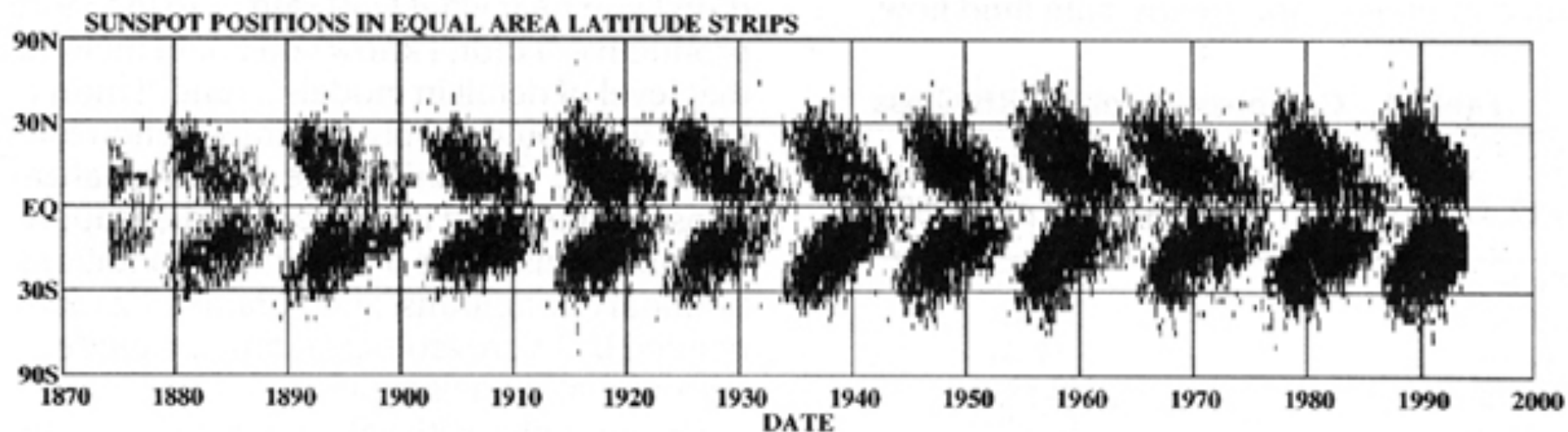
呂凌霄

中央大學太空科學研究所

11-years Solar Cycle Based on Sunspot Number

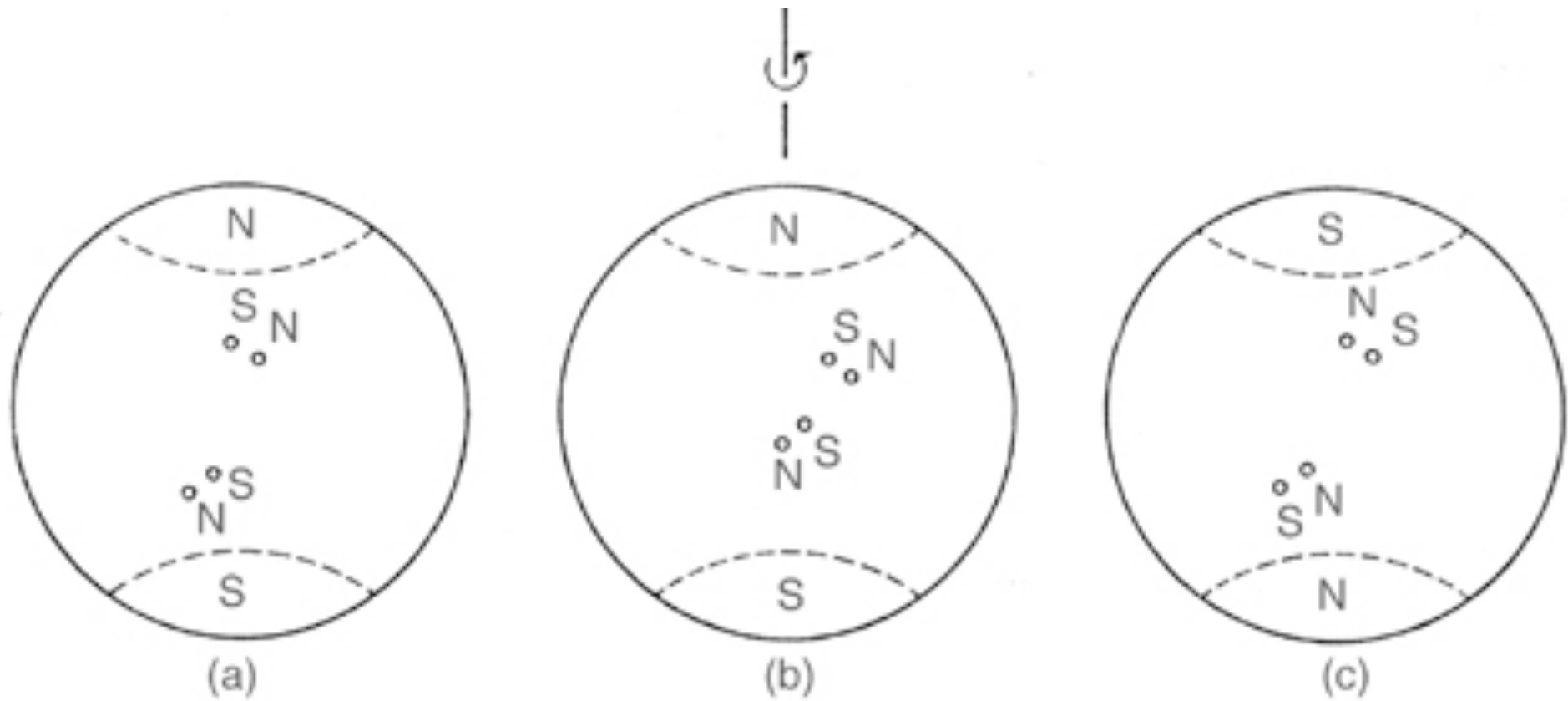


DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



22-years Solar Cycle

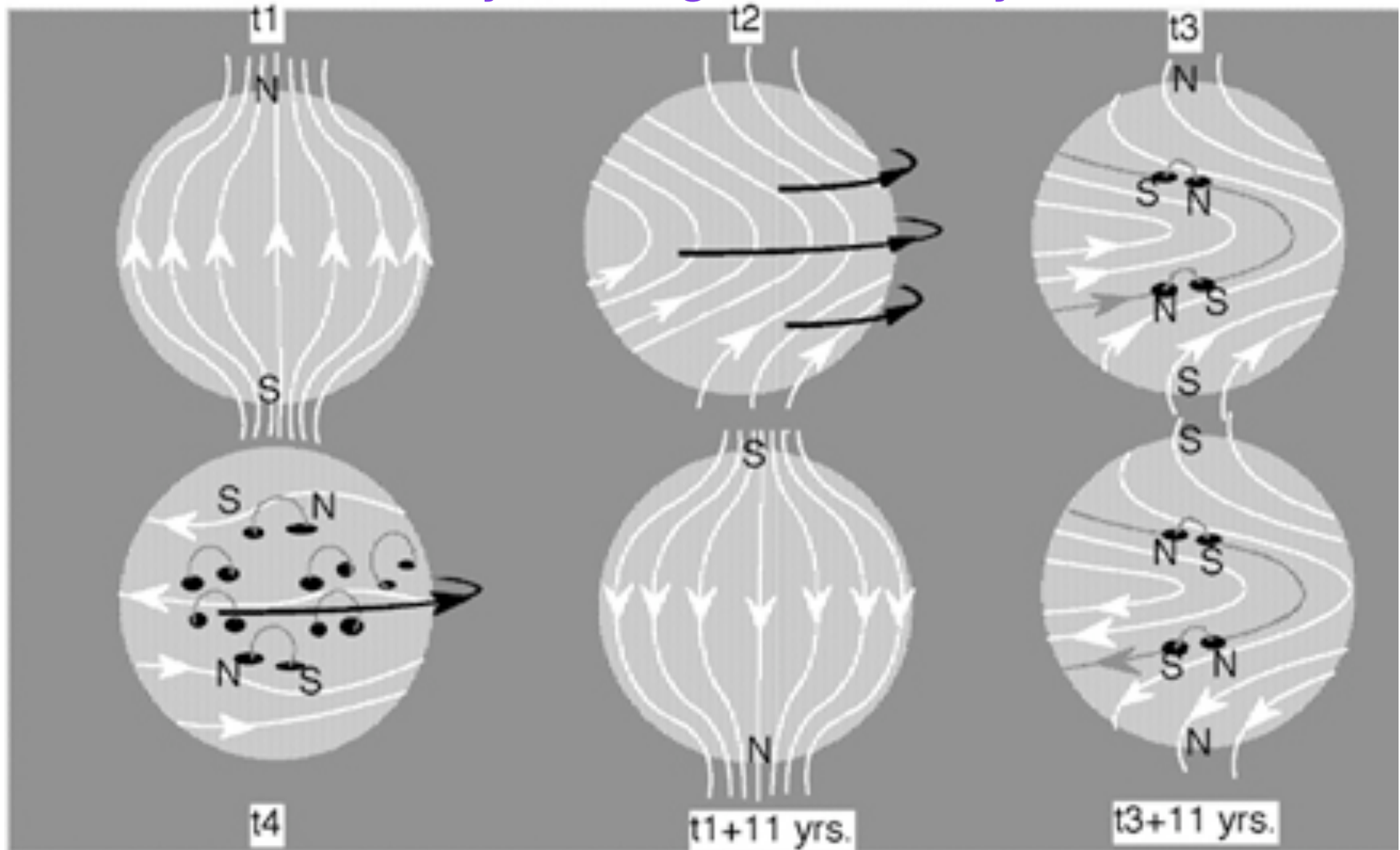
Based on Leading Sunspot Magnetic Polarity



Solar Differential Rotation

--> Solar Dynamo (w-effects, wa-effects)

-->22-years Magnetic Solar Cycle

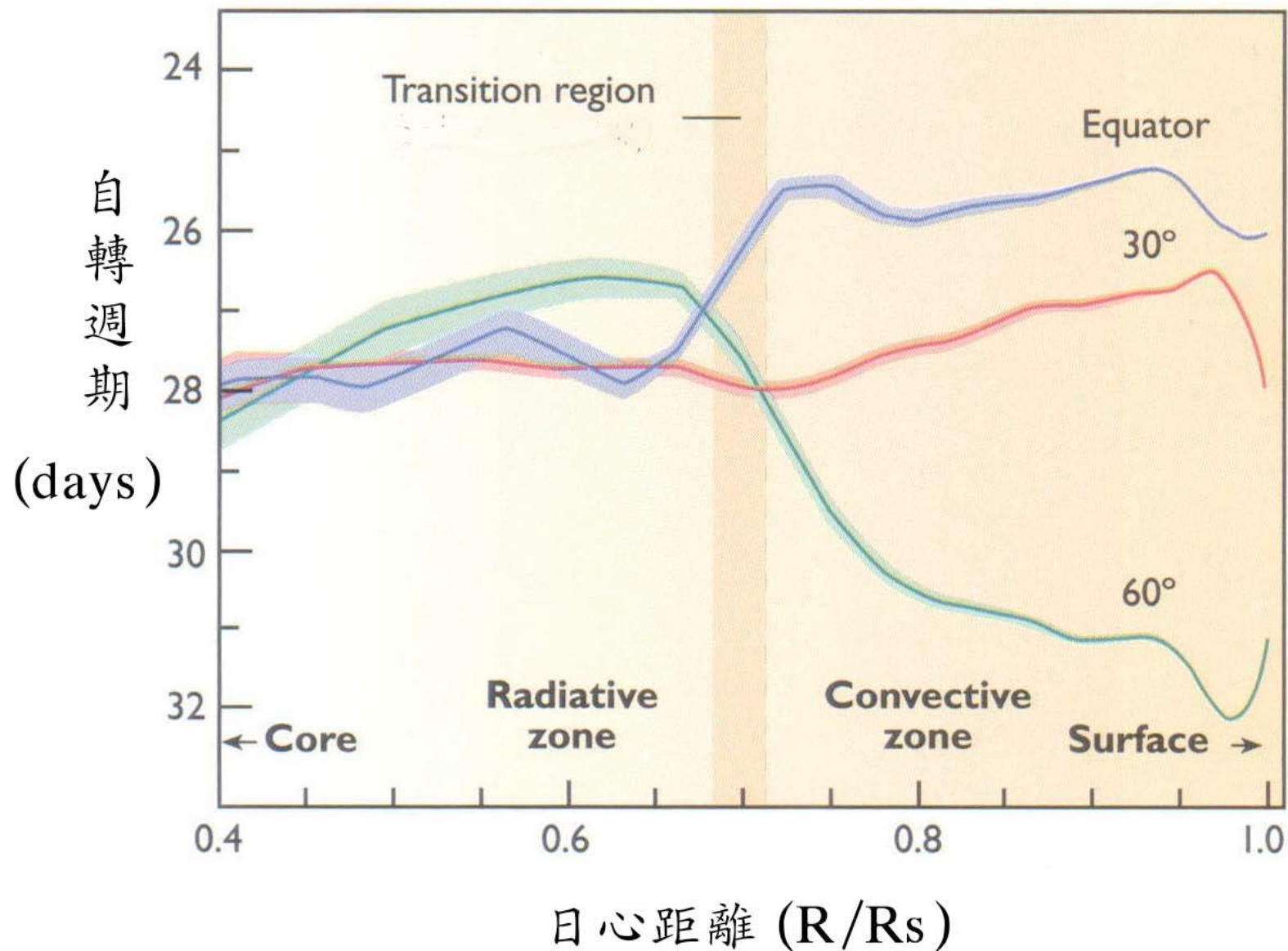


Differential rotation is the major cause of solar cycle!

What is the cause of solar differential rotation?

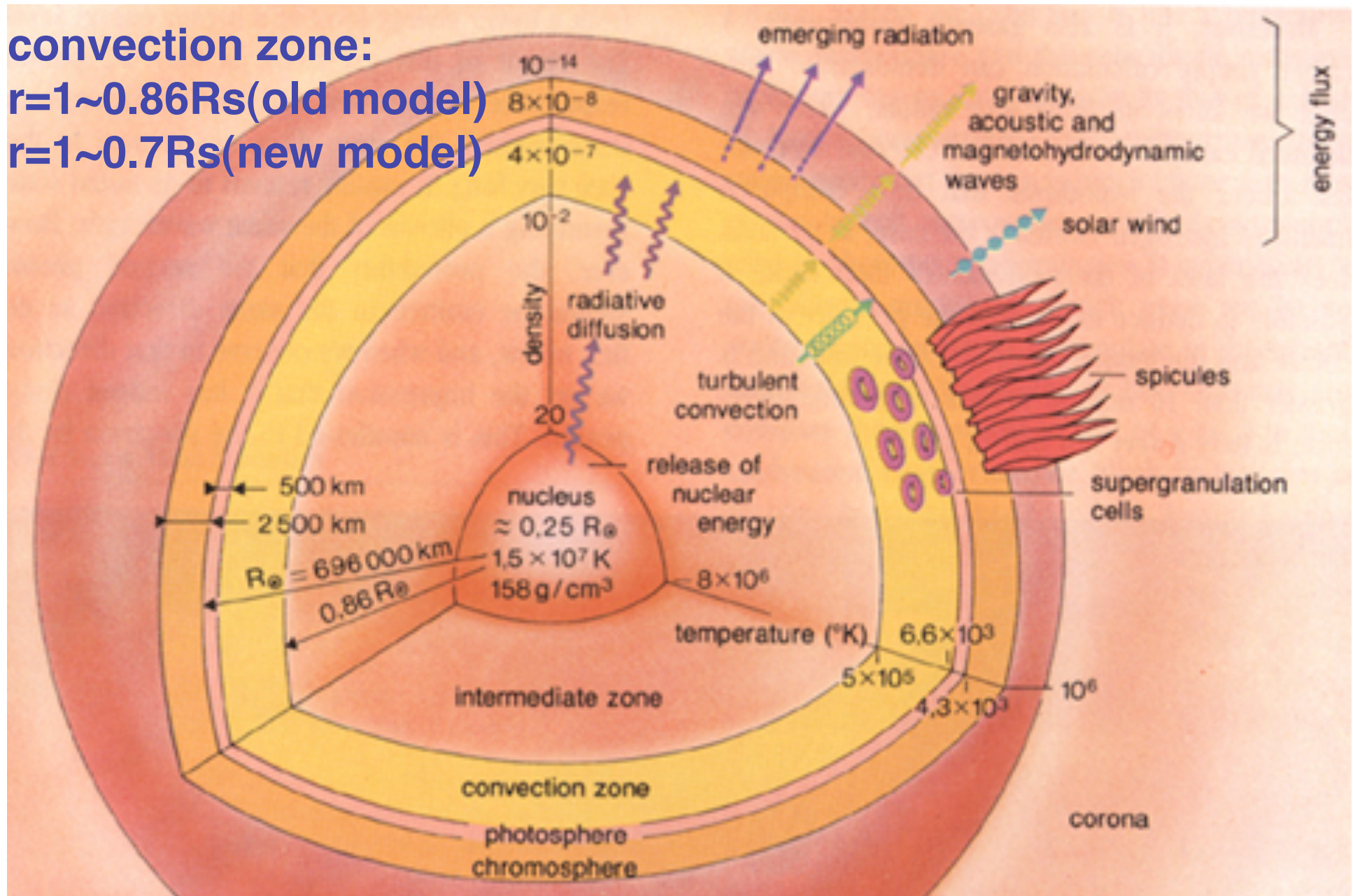
Differential Rotation in Solar Interior

SoHO/MDI Observation



Solar Interior Structures

convection zone:
 $r=1 \sim 0.86 R_{\odot}$ (old model)
 $r=1 \sim 0.7 R_{\odot}$ (new model)



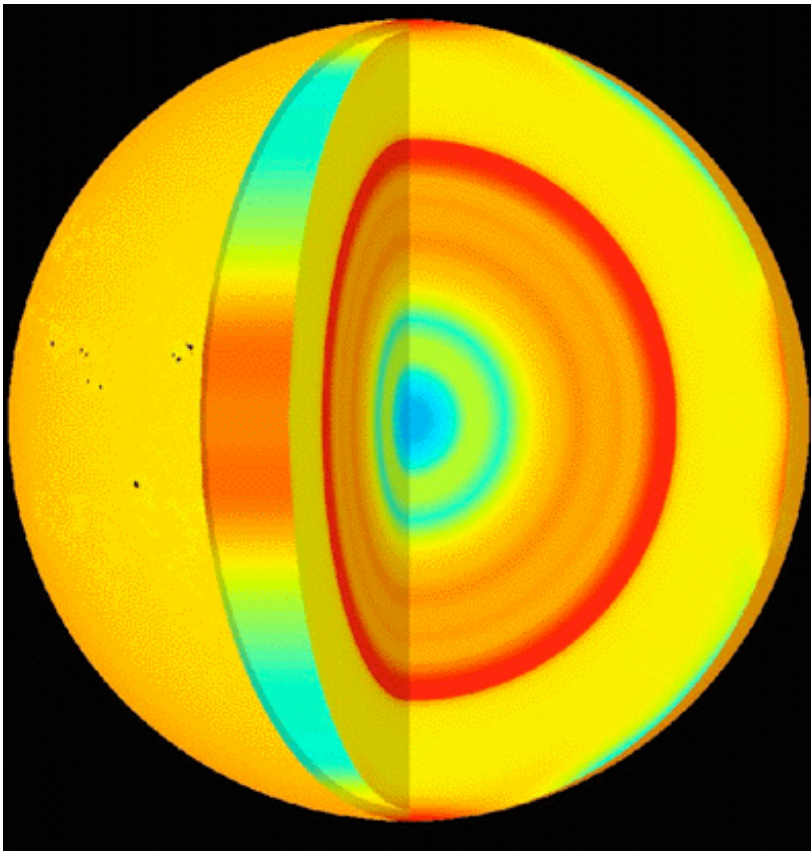
Sound Speed Distribution in Solar Interior

Abnormal plasma sound speed
($c^2 = T/\bar{m}$) distribution in the solar
interior (MDI observation)

$\delta(c^2) > 0$

$\delta(c^2) < 0$

T : plasma temperature
 $\bar{m}$ average mass



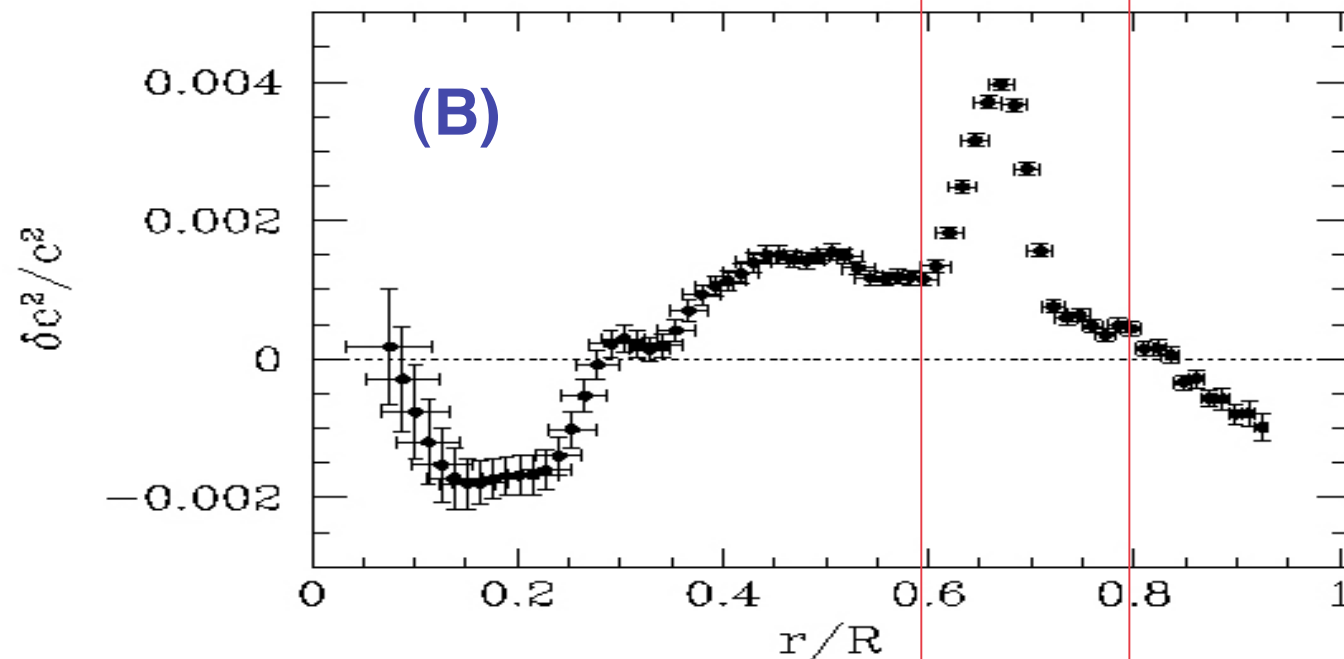
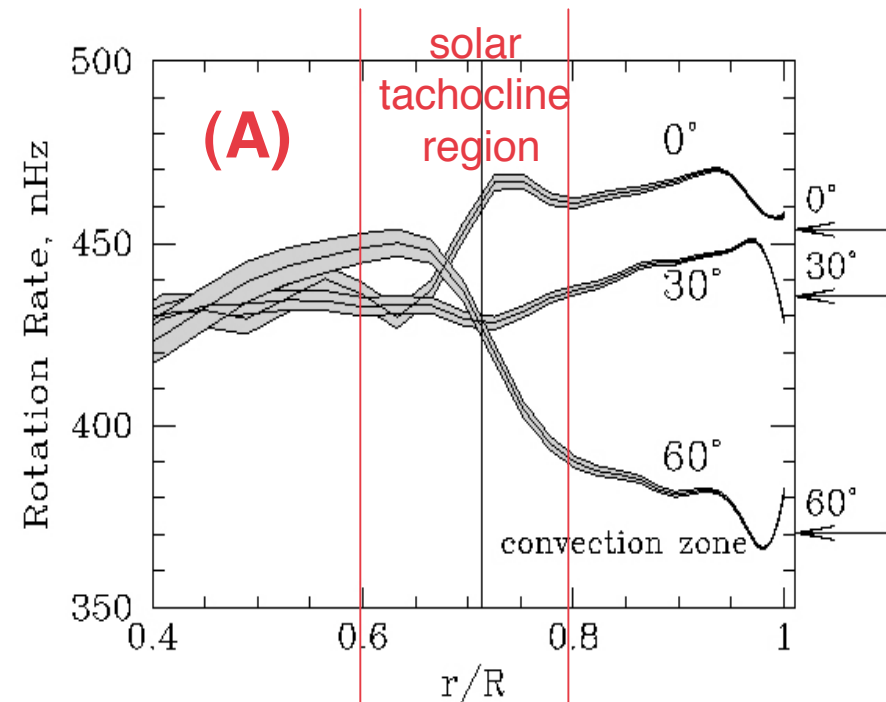
MDI Observations (Kosovichev et al., 1997)

(A) Differential Rotation

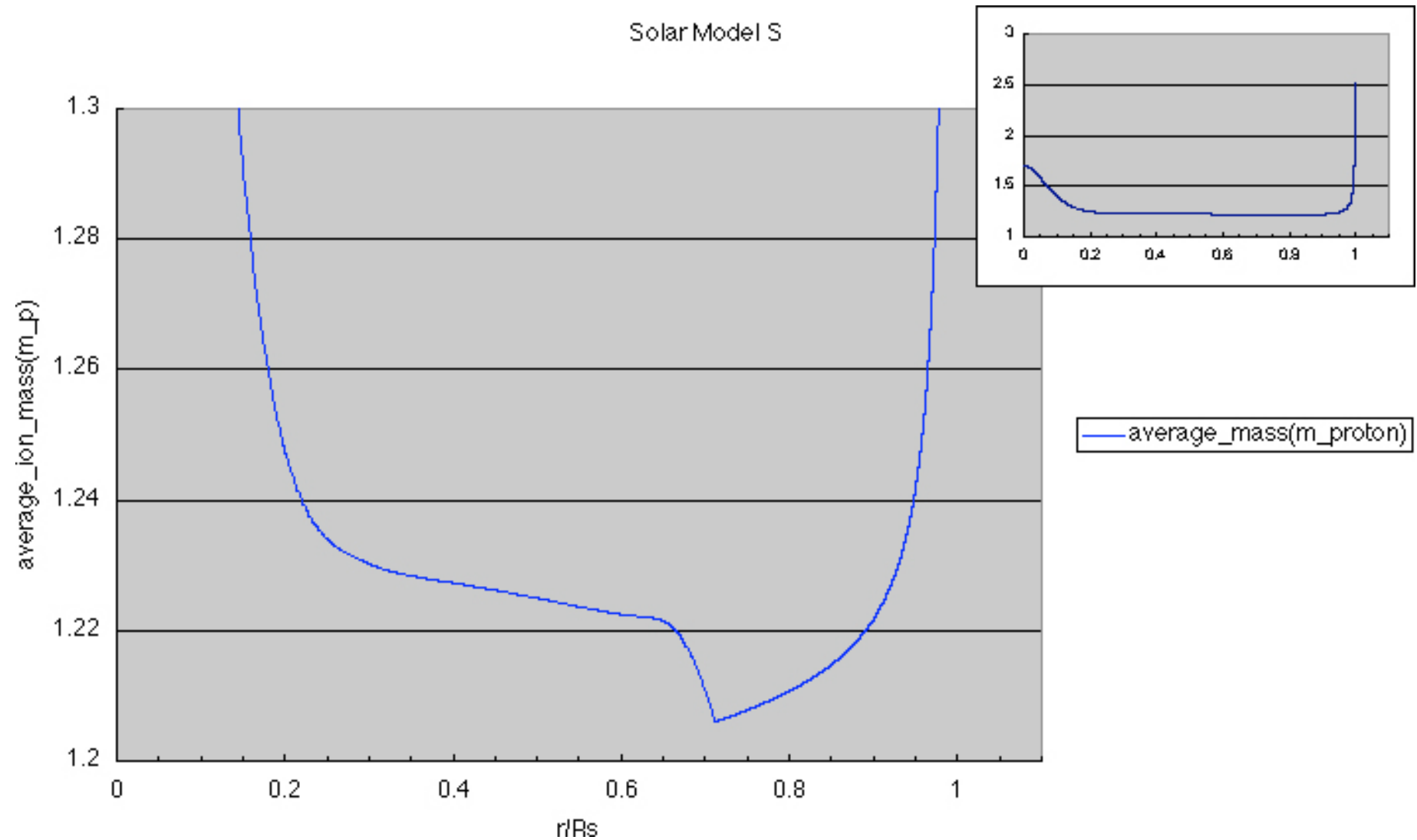
(B) Abnormal Plasma Temperature Distribution

Cause → Result

(A) → *Turbulence Heating* → (B)



Average Mass Distribution in Standard Solar Model



(C) Average Mass m Distribution in Standard Solar Model

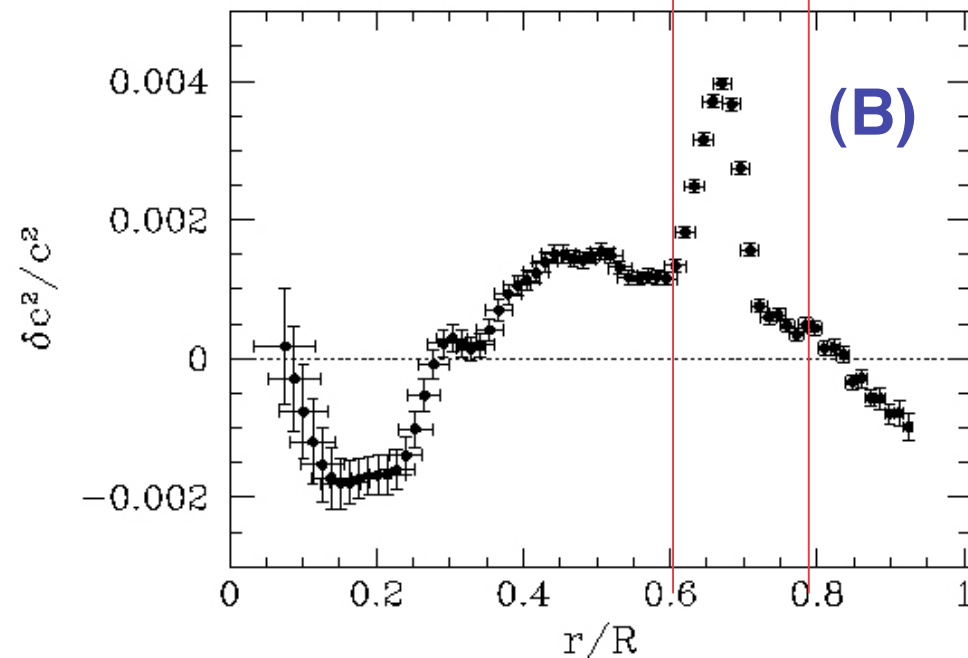
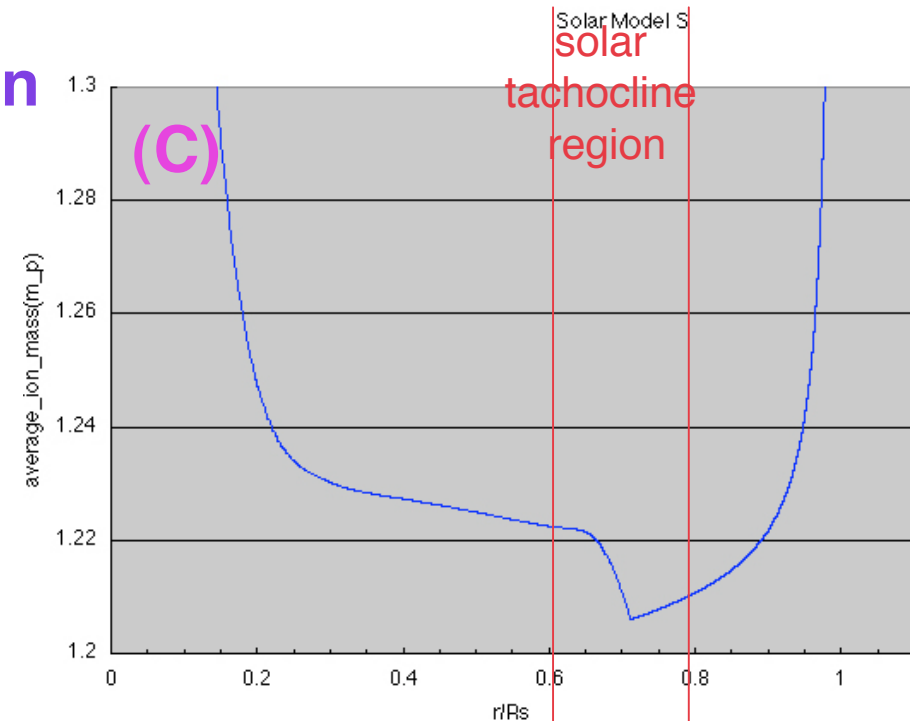
Cause → Result

Under estimate mixing length

→ Under estimate m

→ Over estimate $c^2 = T/m$

→ (B) Abnormal Plasma
Temperature Distribution



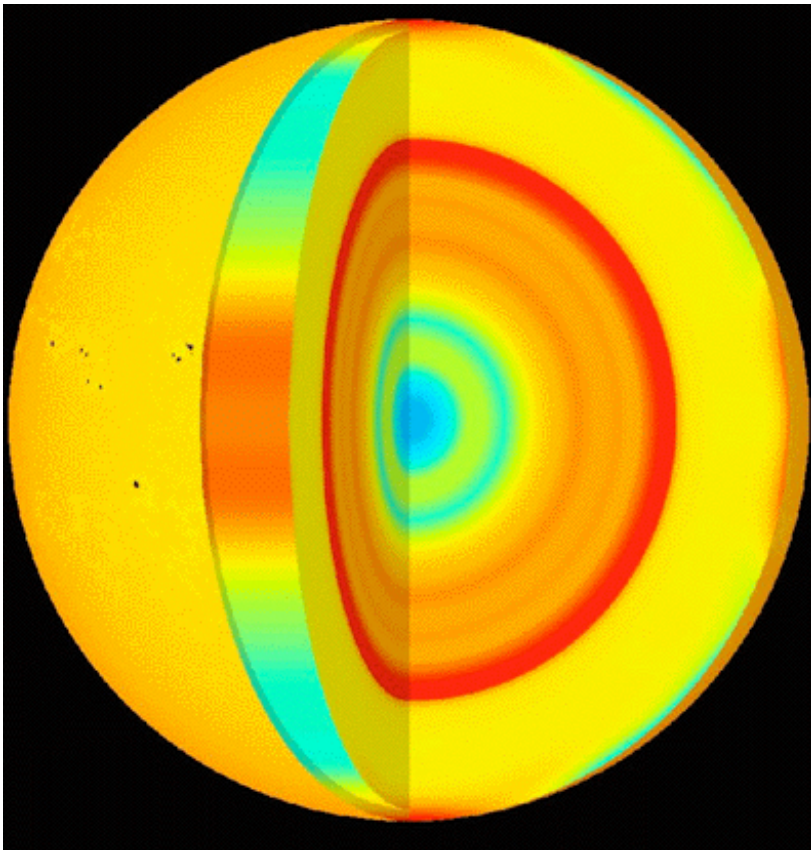
Sound Speed Distribution in Solar Interior

Abnormal plasma sound speed
($c^2 = T/\bar{m}$) distribution in the solar
interior (MDI observation)

$\delta(c^2) > 0$

$\delta(c^2) < 0$

T : plasma temperature
 $\bar{m}$ average mass



Our Explanation

We assume that a *normal* state means a hydrodynamic equilibrium state.

$$-\nabla p + \rho \mathbf{g} = 0 \quad \text{or} \quad -\frac{dp}{dr} = \rho g$$

Then, *abnormal* temperature profile in the solar tachocline region

indicates that $-\frac{dp}{dr} > \rho g$

To reach a steady state, there must be an *inward force* to make force balance in the r-direction.

We propose that

(1) abnormal temperature only occur in the middle and lower latitude region,

(2) the inward force is a $\mathbf{J} \times \mathbf{B}$ force, such that

$$\left(-\frac{dp}{dr} - \rho g \right) \hat{r} + \mathbf{J} \times \mathbf{B} = 0 \quad \text{or} \quad \mathbf{J} = \left(-\frac{dp}{dr} - \rho g \right) \hat{r} \times \frac{\mathbf{B}}{B^2}$$

(3) absence of such inward force ($\mathbf{J} \times \mathbf{B}$) in the polar coronal hole region is because that $\mathbf{B} = B_r \hat{r}$ in the polar coronal hole region.

Our Explanation on Abnormal c^2 Distribution

Abnormal plasma sound speed
($c^2=T/\bar{m}$) distribution in the solar
interior (MDI observation)

$\delta(c^2)>0$

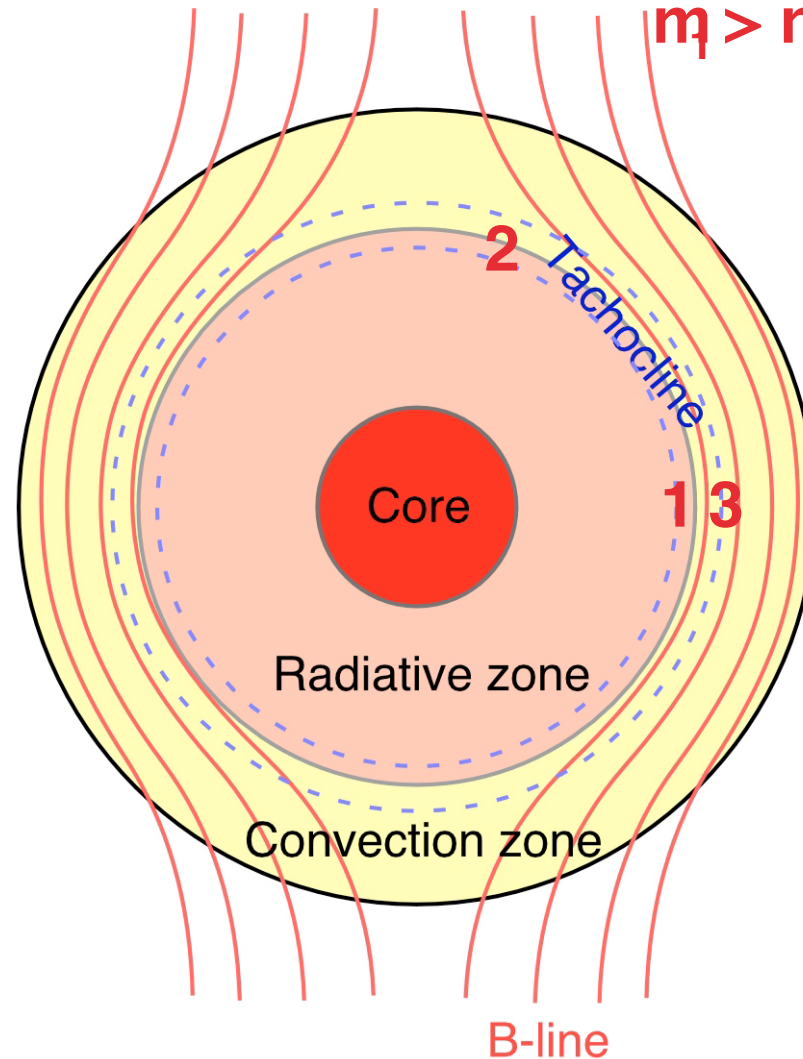
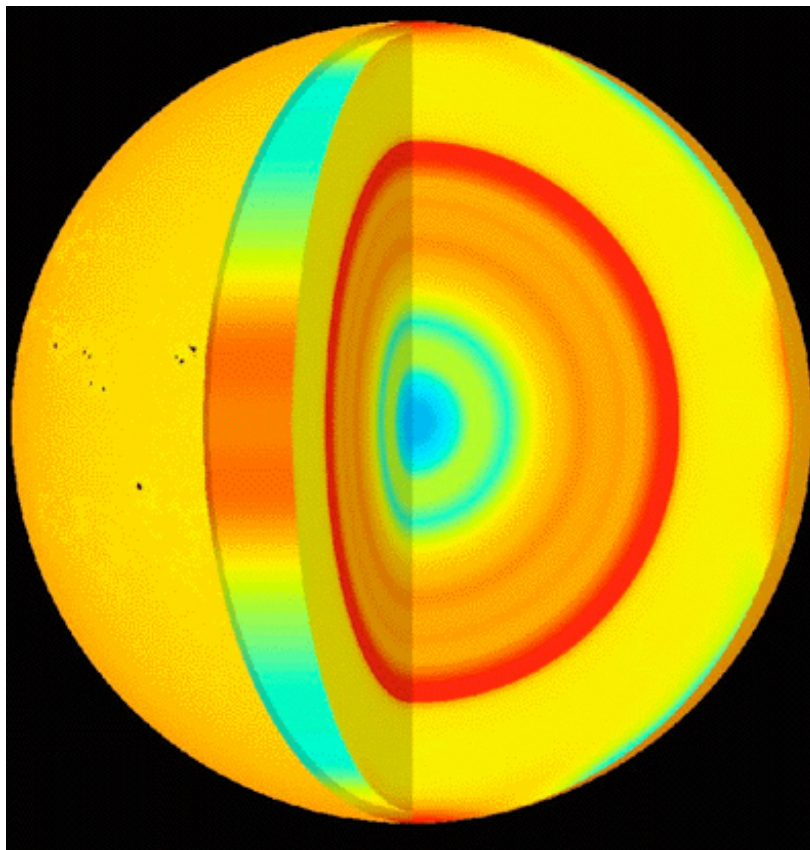
$\delta(c^2)<0$

Plasma Temperature

$$T_1 > T_2 > T_3$$

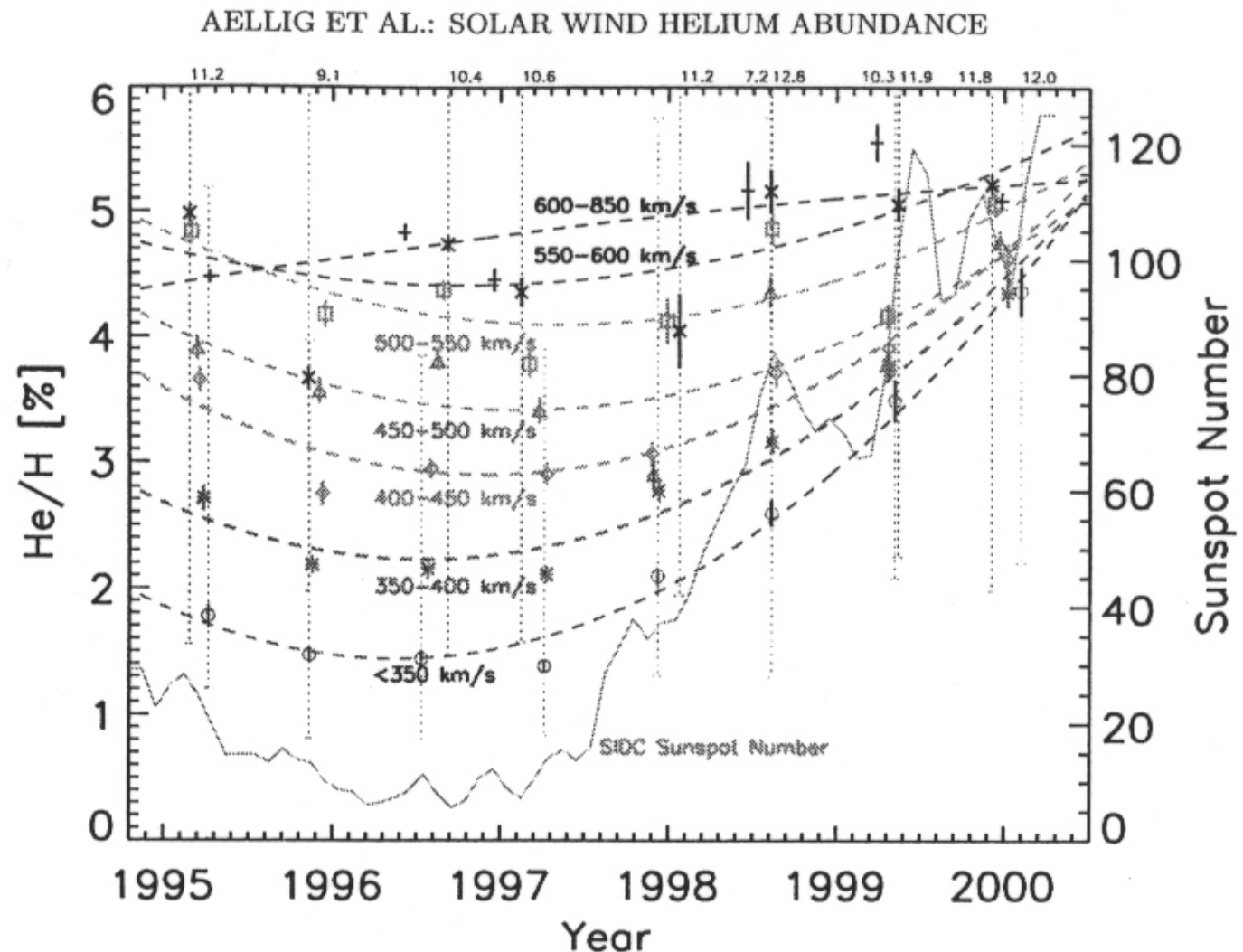
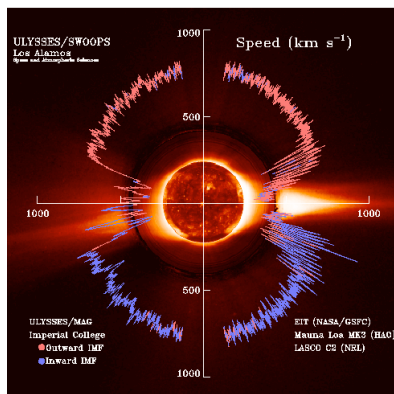
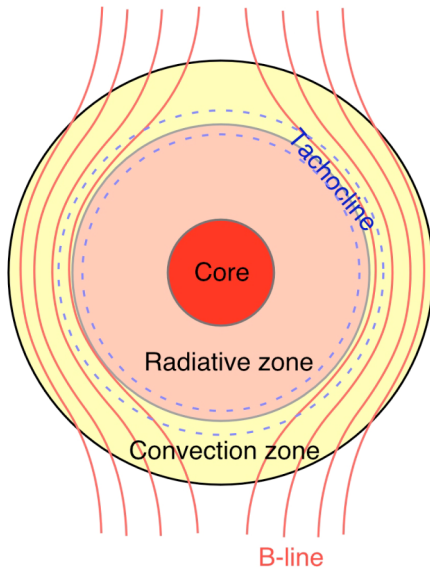
Average Mass:

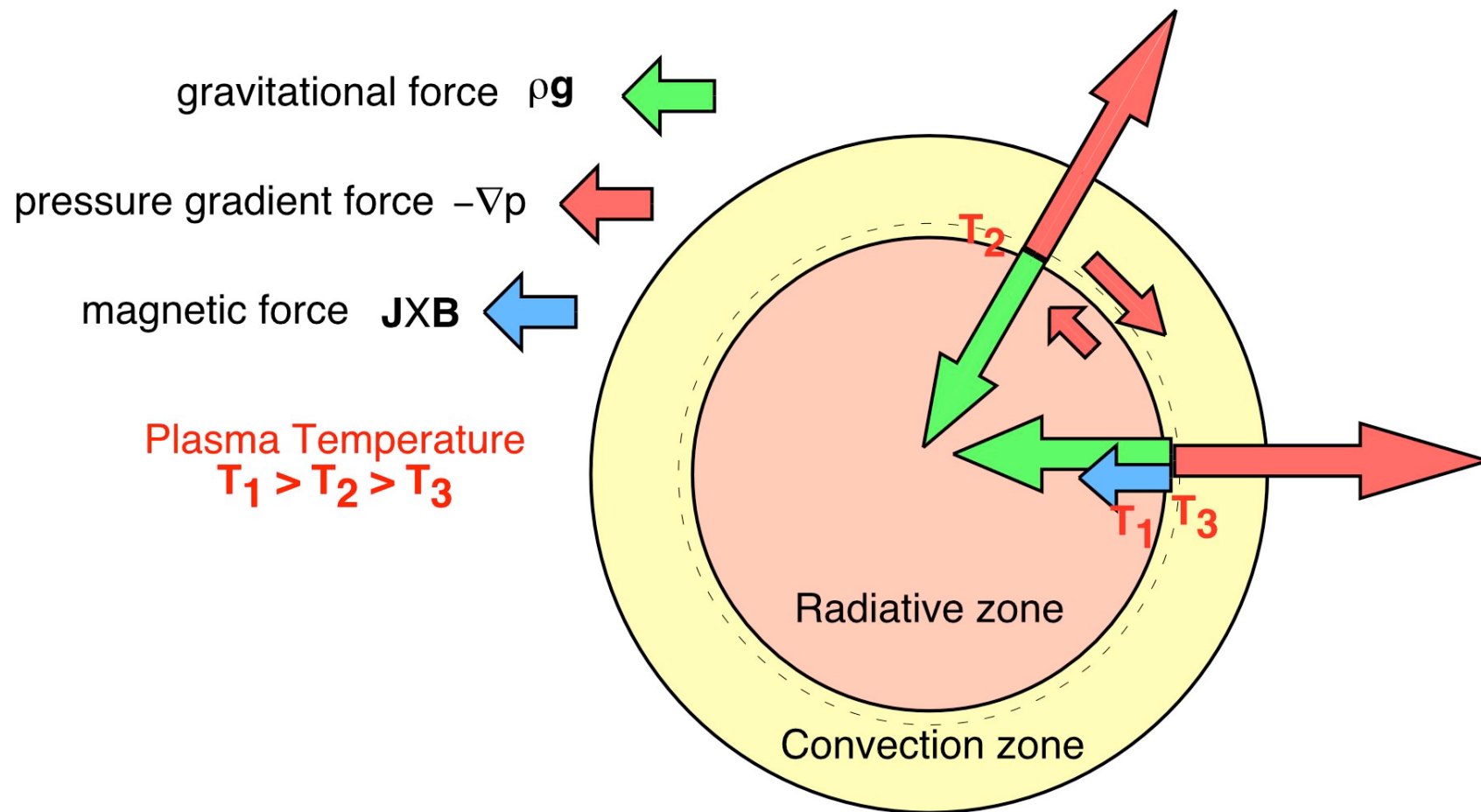
$$\bar{m}_1 > \bar{m}_2 > \bar{m}_3$$



Evidence on Smaller mixing length in the low latitude region Large mixing length in the high latitude region

Average mass
 $m_1 > m_2 > m_3$





Zonal flow can lead to horizontal force balance in a rotational star or planet with latitudinal temperature gradient

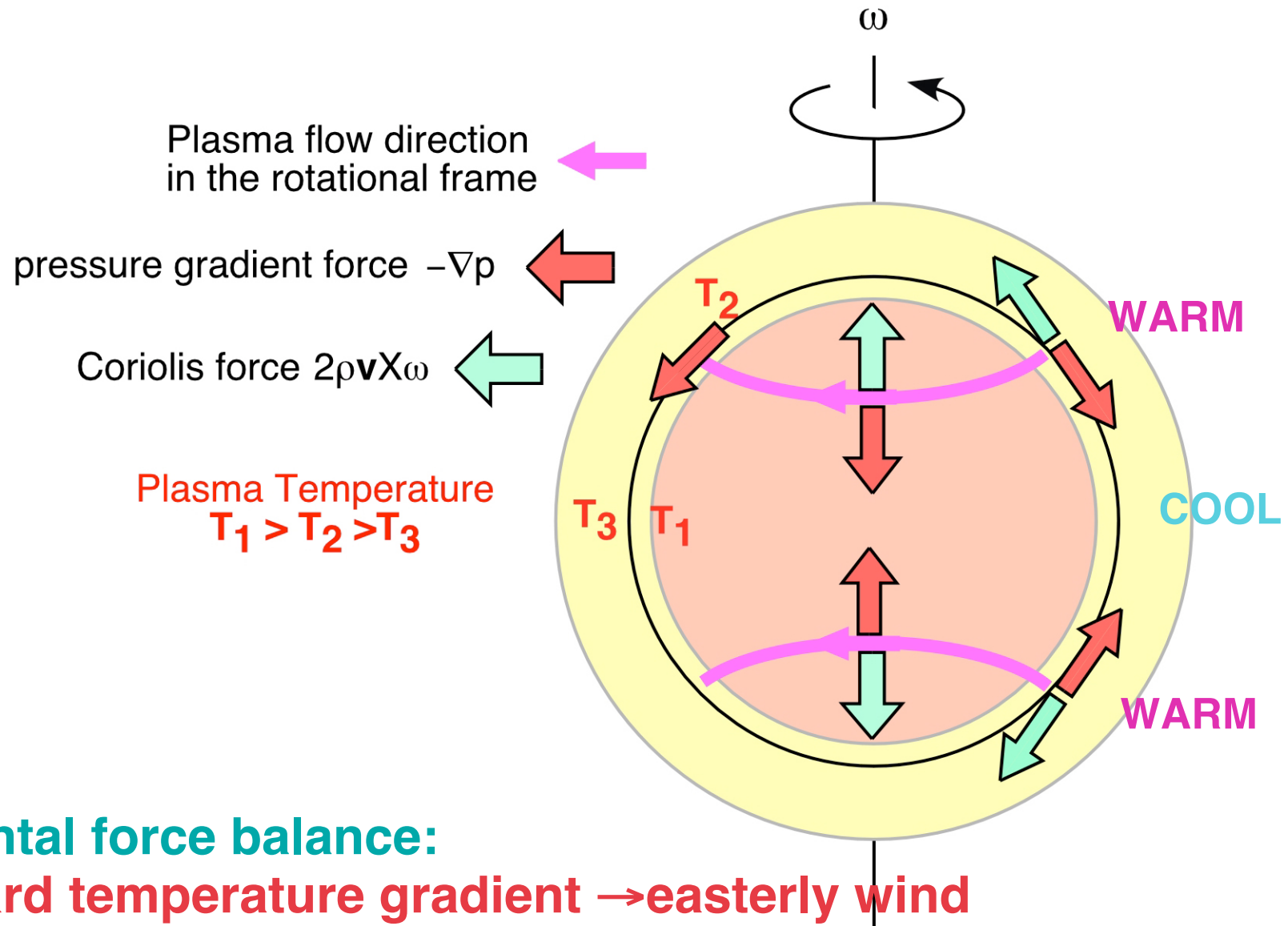
Upper-half tachocline: (in solar convection zone)

Poleward temperature gradient $\rightarrow$ easterly wind

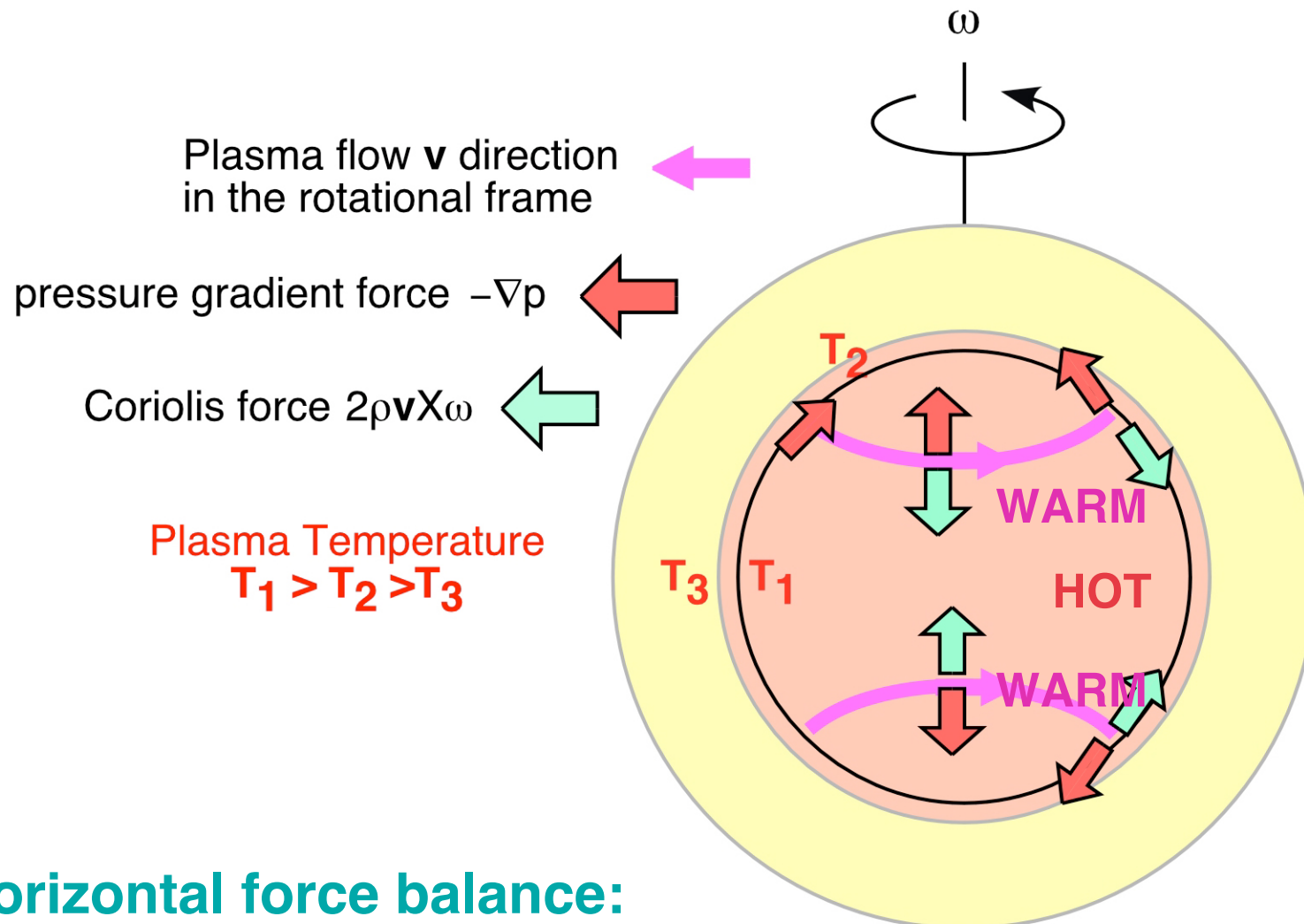
Lower-half tachocline: (in solar radiative zone)

Equatorward temperature gradient $\rightarrow$ westerly wind

Easterly wind in the upper-half solar tachocline region (in solar convection zone).

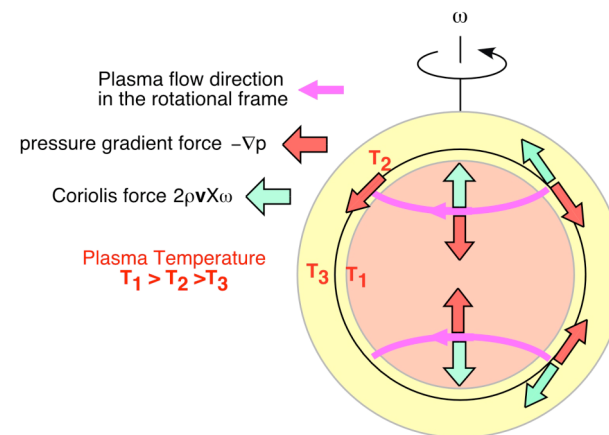
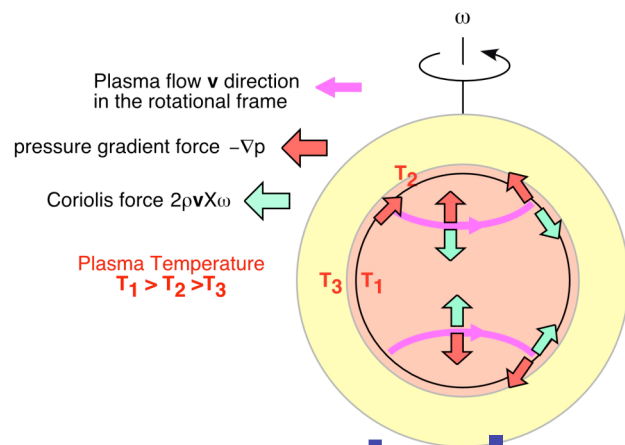


Westerly wind in the lower-half solar tachocline region (in solar radiative zone)



Horizontal force balance:

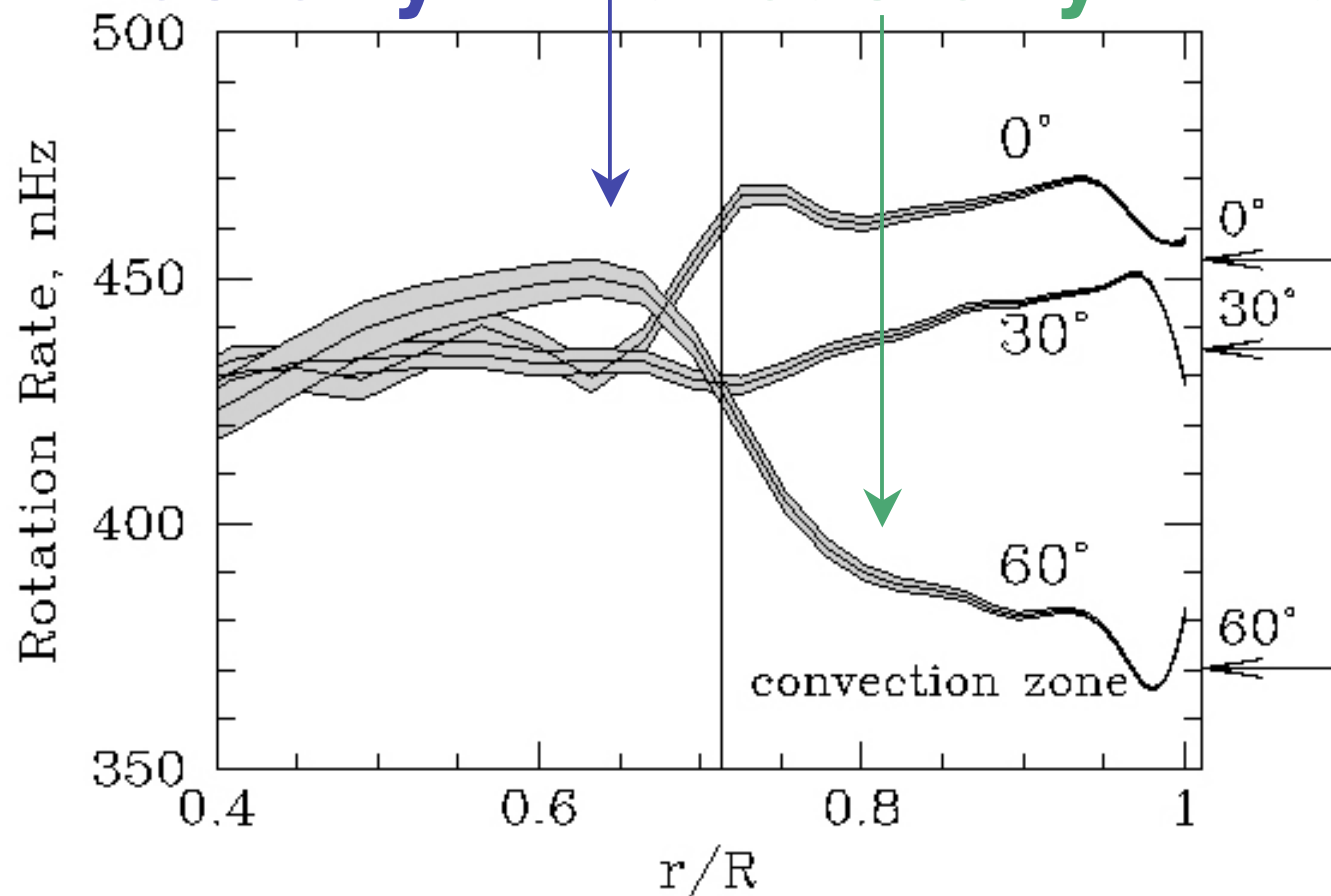
Equatorward temperature gradient $\rightarrow$ westerly wind



~60 degrees

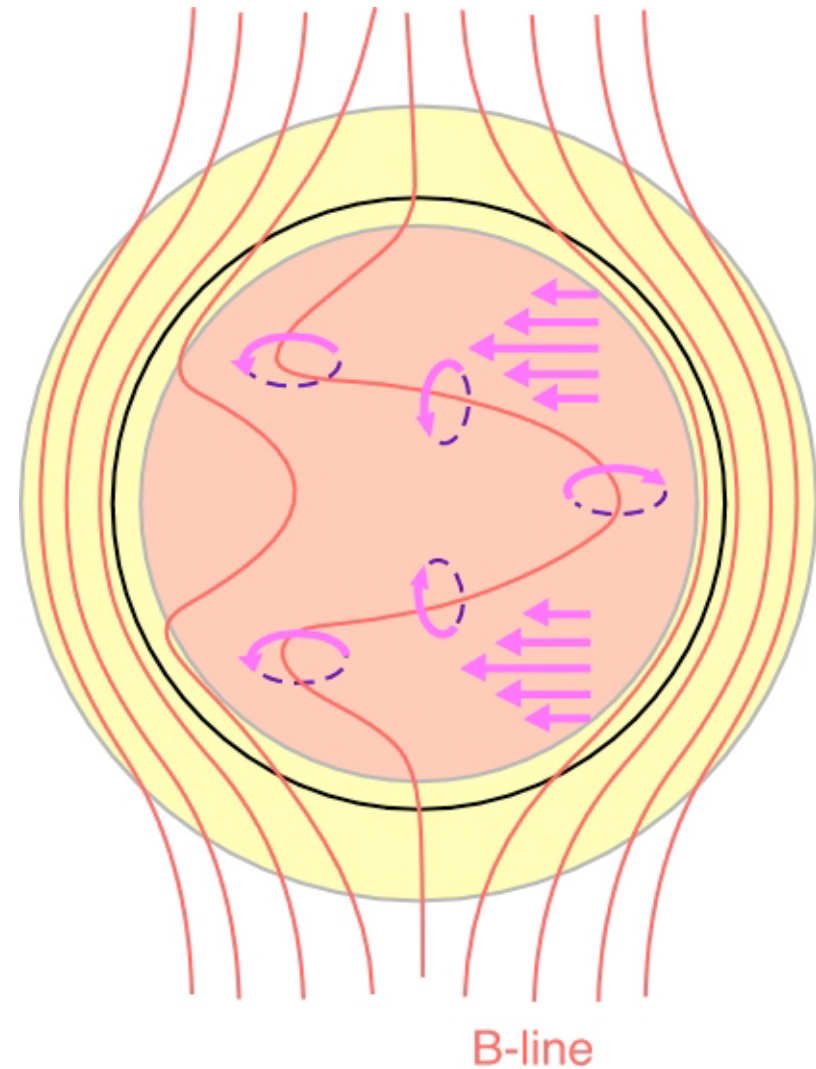
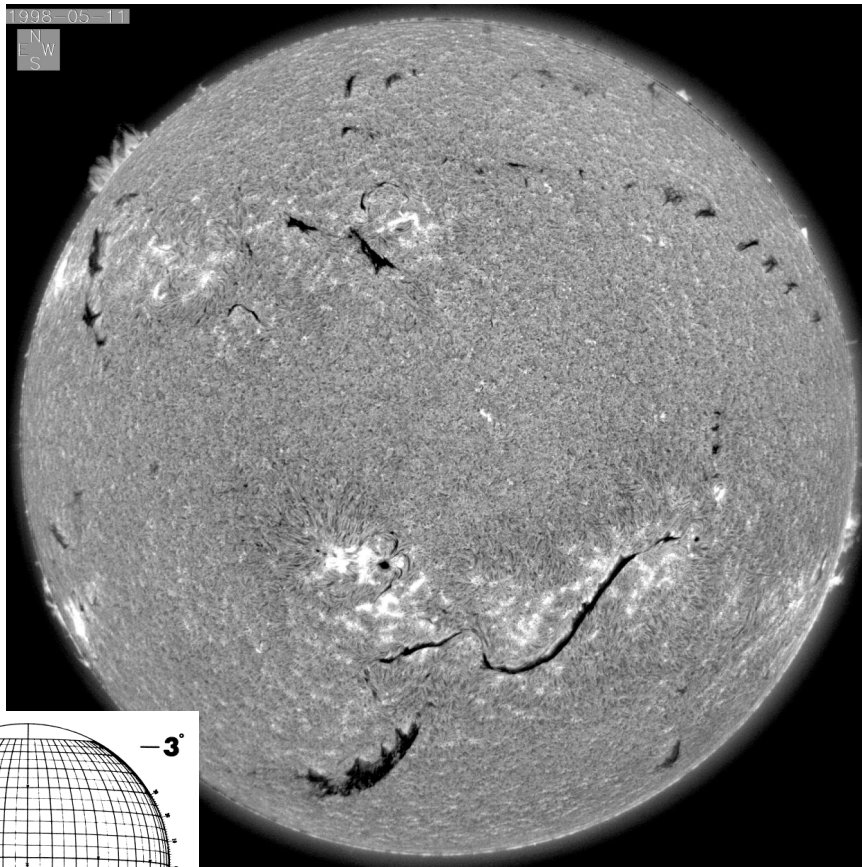
westerly wind

easterly wind

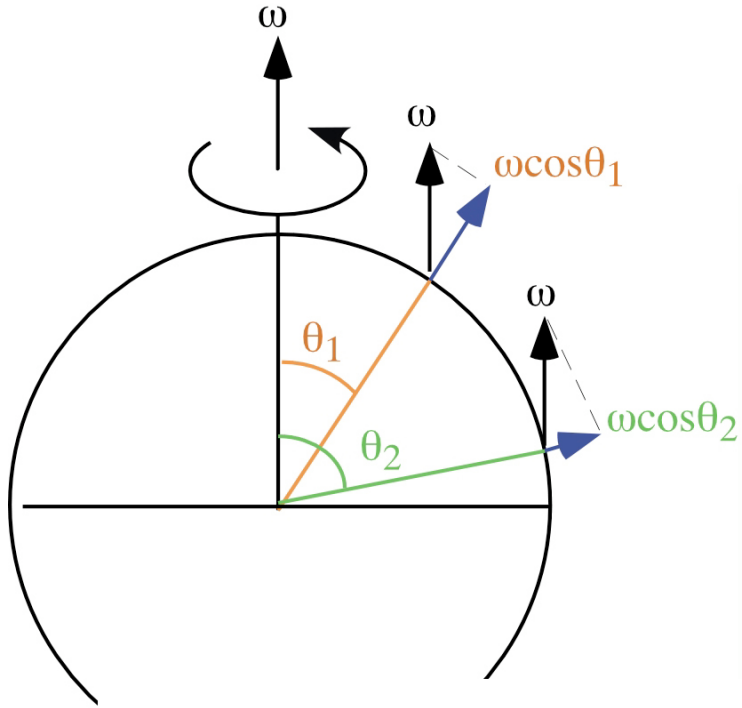


Initial formation of solar filaments' magnetic field

BBO H α observation
of solar filaments 05/11/1998



Zonal Flow in Rotational System --> Rossby wave



momentum equation

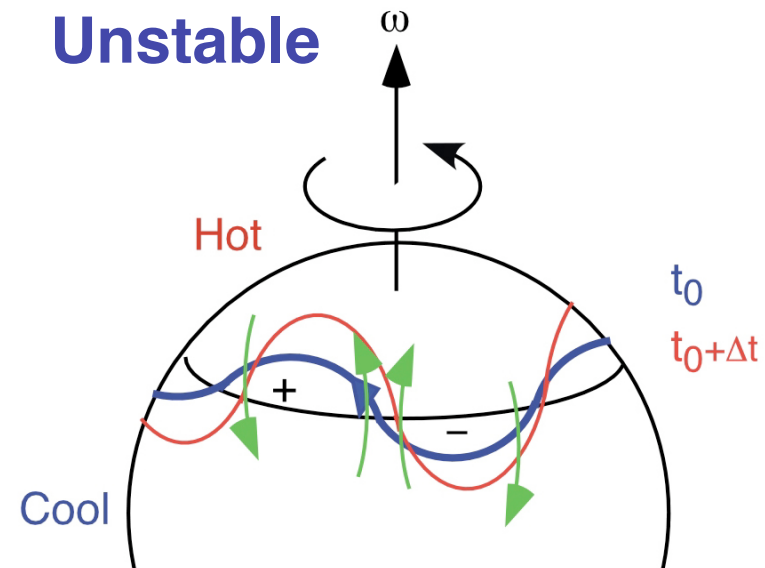
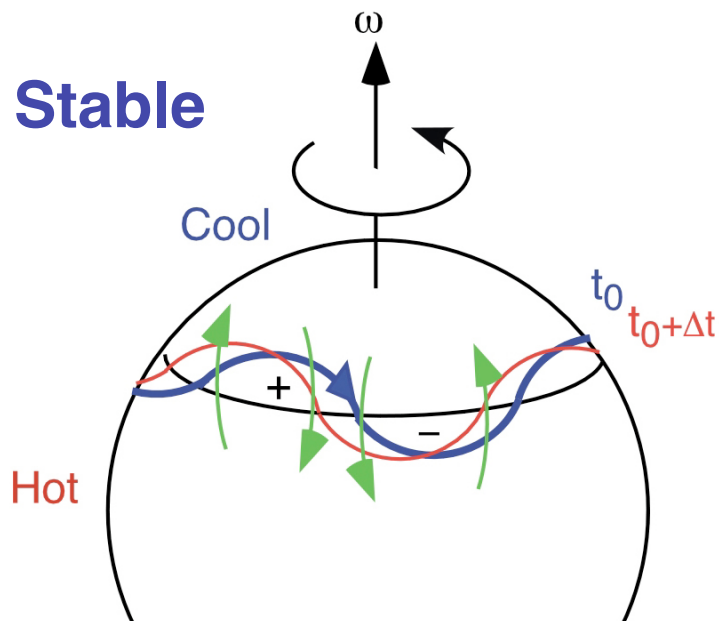
$$\rho(\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V}) = -\nabla p - 2\rho \boldsymbol{\omega} \times \mathbf{V} + \rho \mathbf{g} - \rho \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) + \mathbf{J} \times \mathbf{B}$$

or

$$\frac{\partial \mathbf{V}}{\partial t} = -\frac{\nabla p}{\rho} - \nabla\left(\frac{V^2}{2}\right) - [(2\boldsymbol{\omega} + \boldsymbol{\Omega}) \times \mathbf{V}] + \mathbf{g}_e + \frac{\mathbf{J} \times \mathbf{B}}{\rho}$$

where

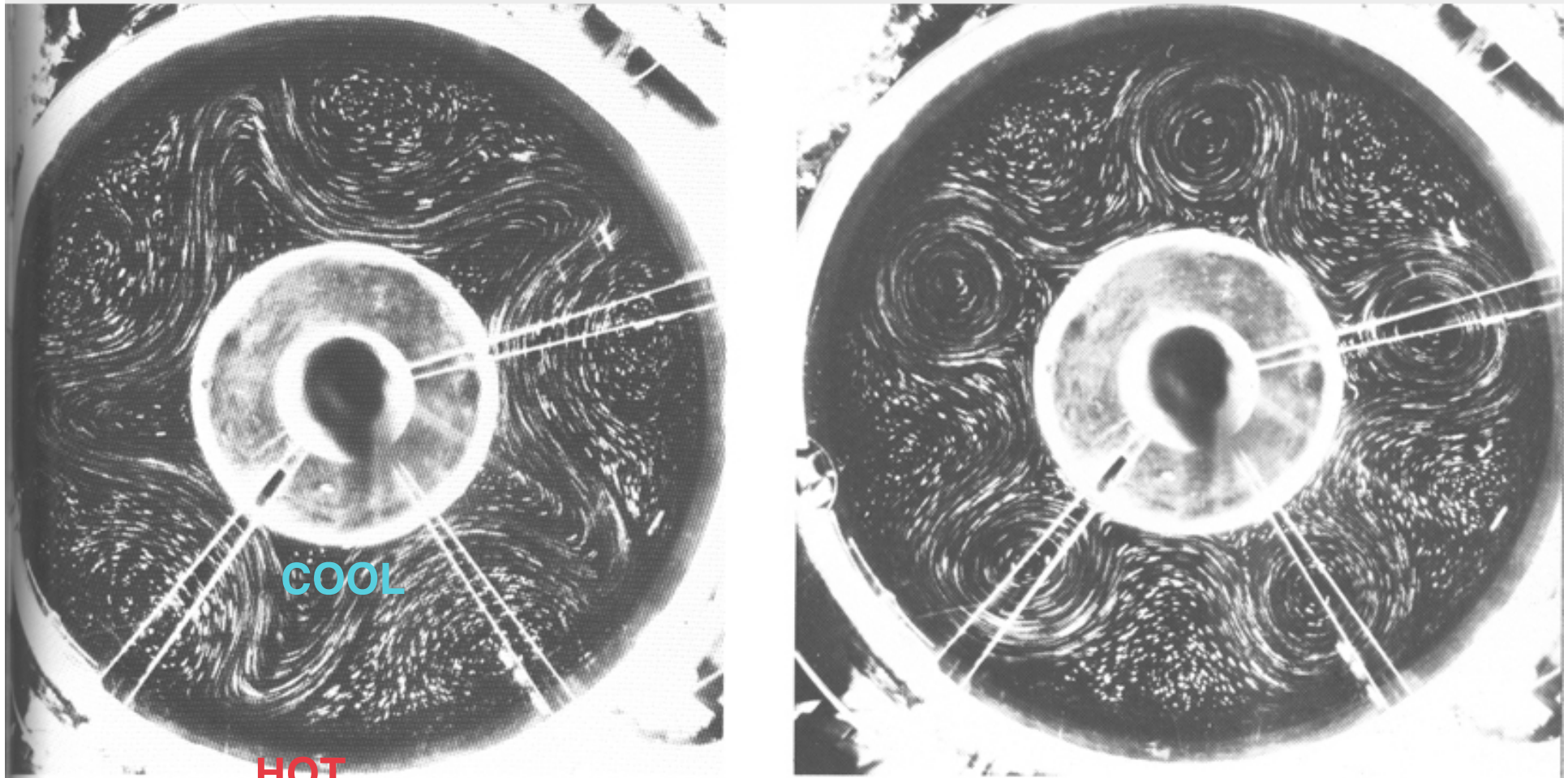
$$\mathbf{\Omega} = \nabla \times \mathbf{V}, \quad \mathbf{g}_e = \mathbf{g} - \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})$$



Laboratory simulation of general circulation (1)

Holton (1992): This is not a Rossby wave, because w is uniform in this case.

Fultz & Spence, 1967.

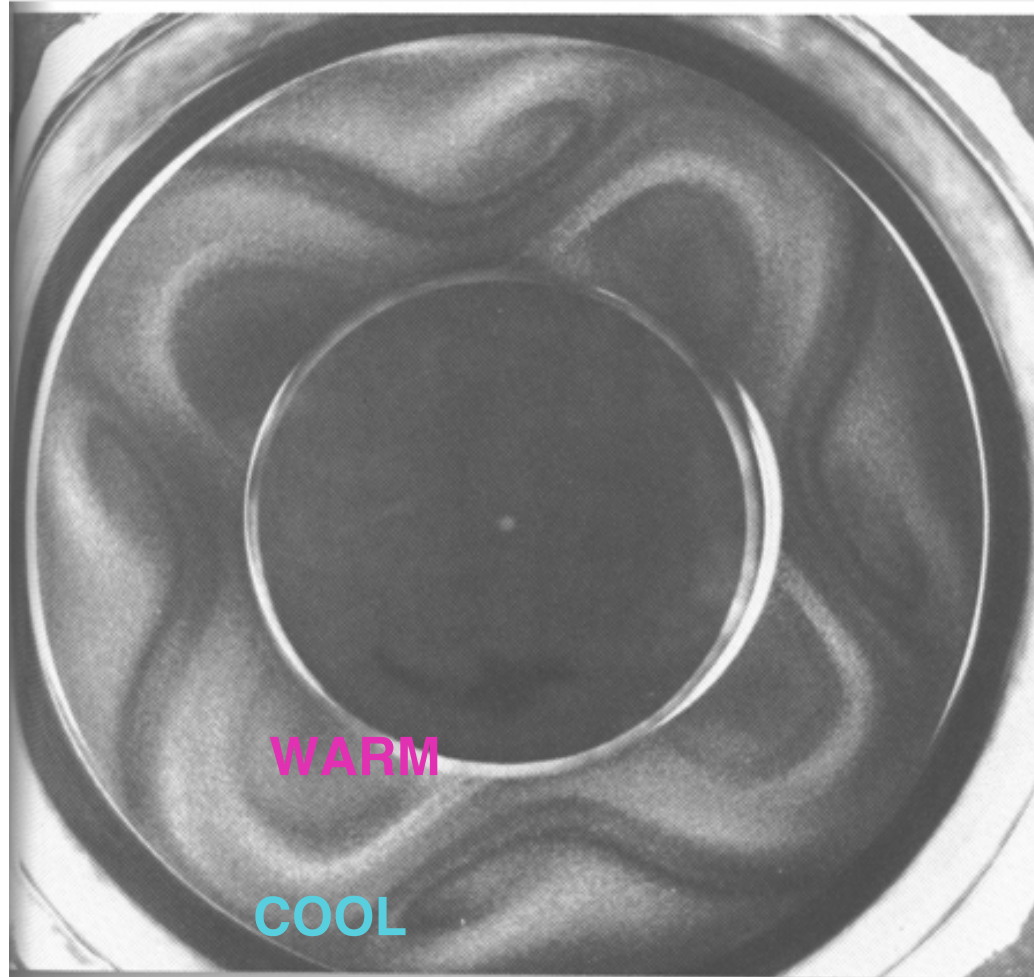


Heating: 80W at the rim and 15 W from a base ring.

Laboratory simulation of general circulation (2)

Holton (1992): This is not a Rossby wave, because ω is uniform in this case.

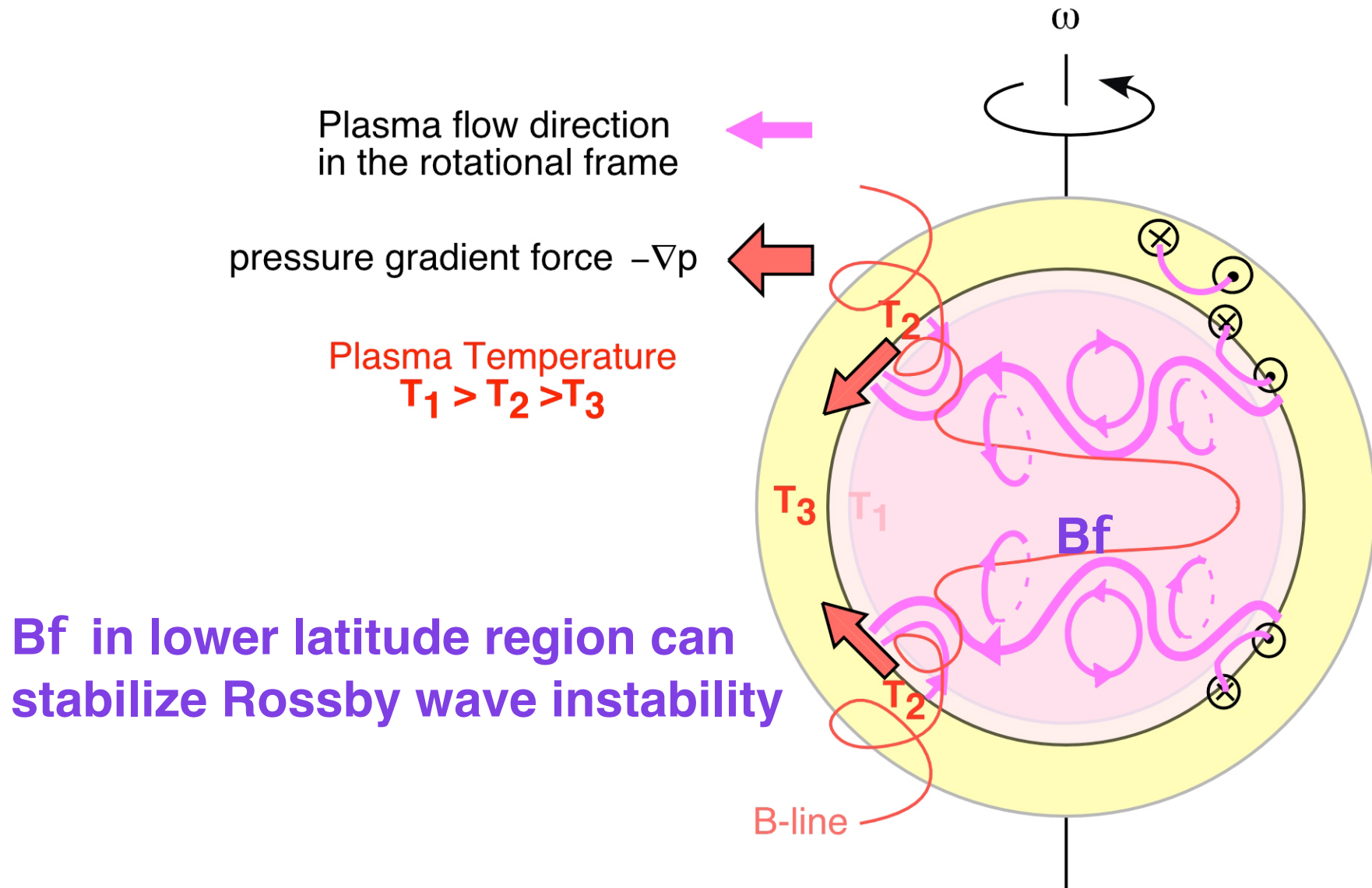
Koschmieder, 1972.



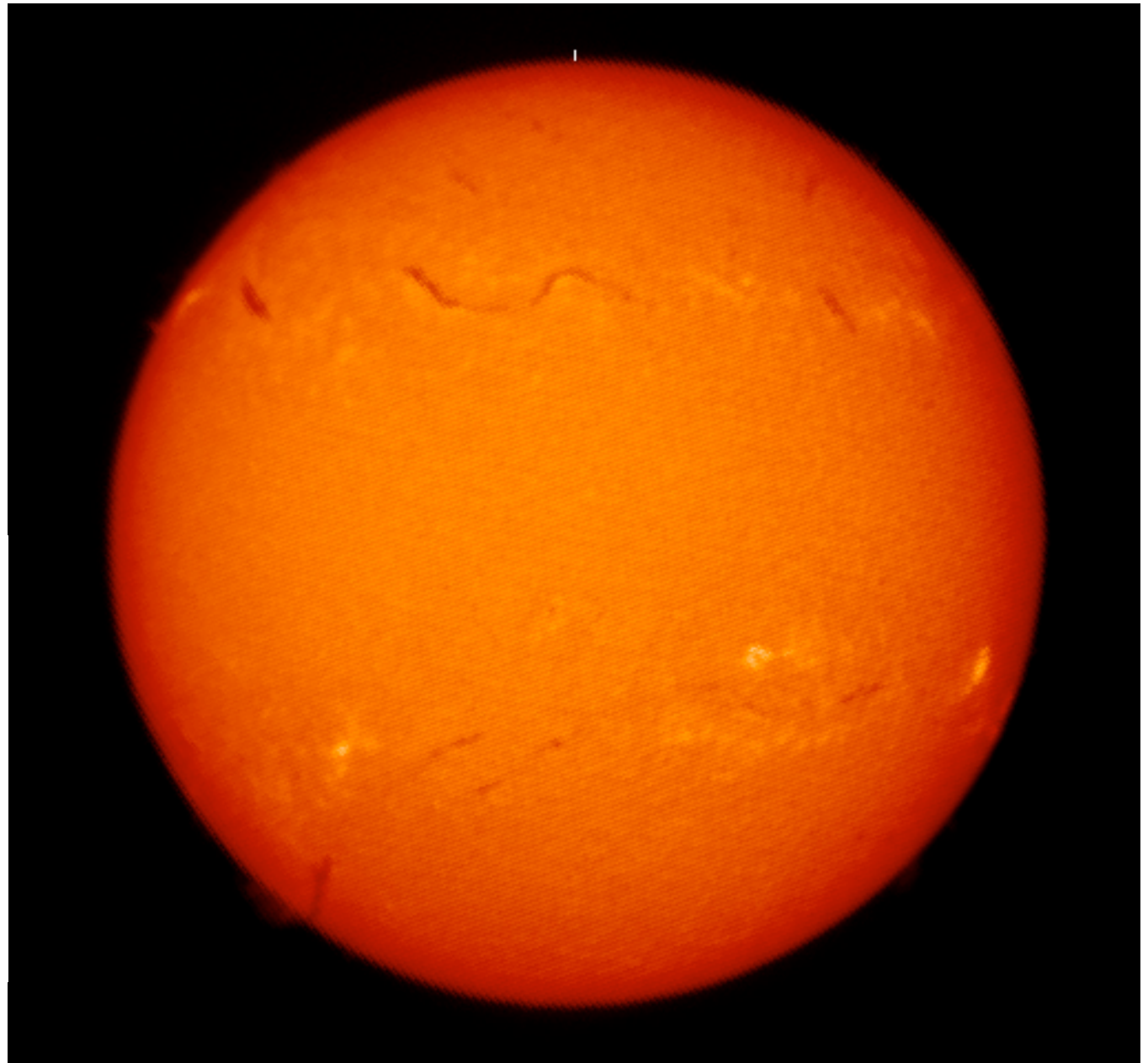
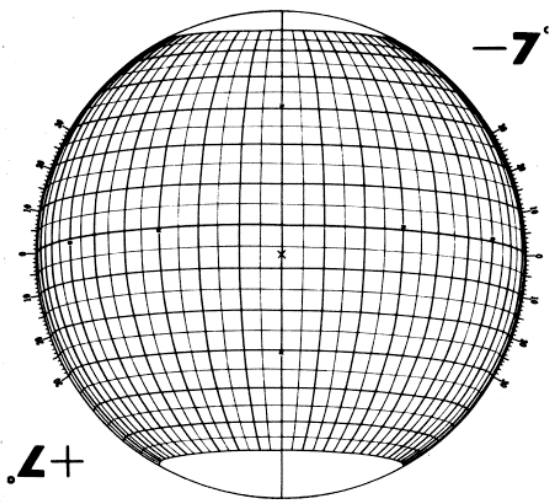
e.g., Solar
convection
zone

The core is 15 degrees warmer than the rim.

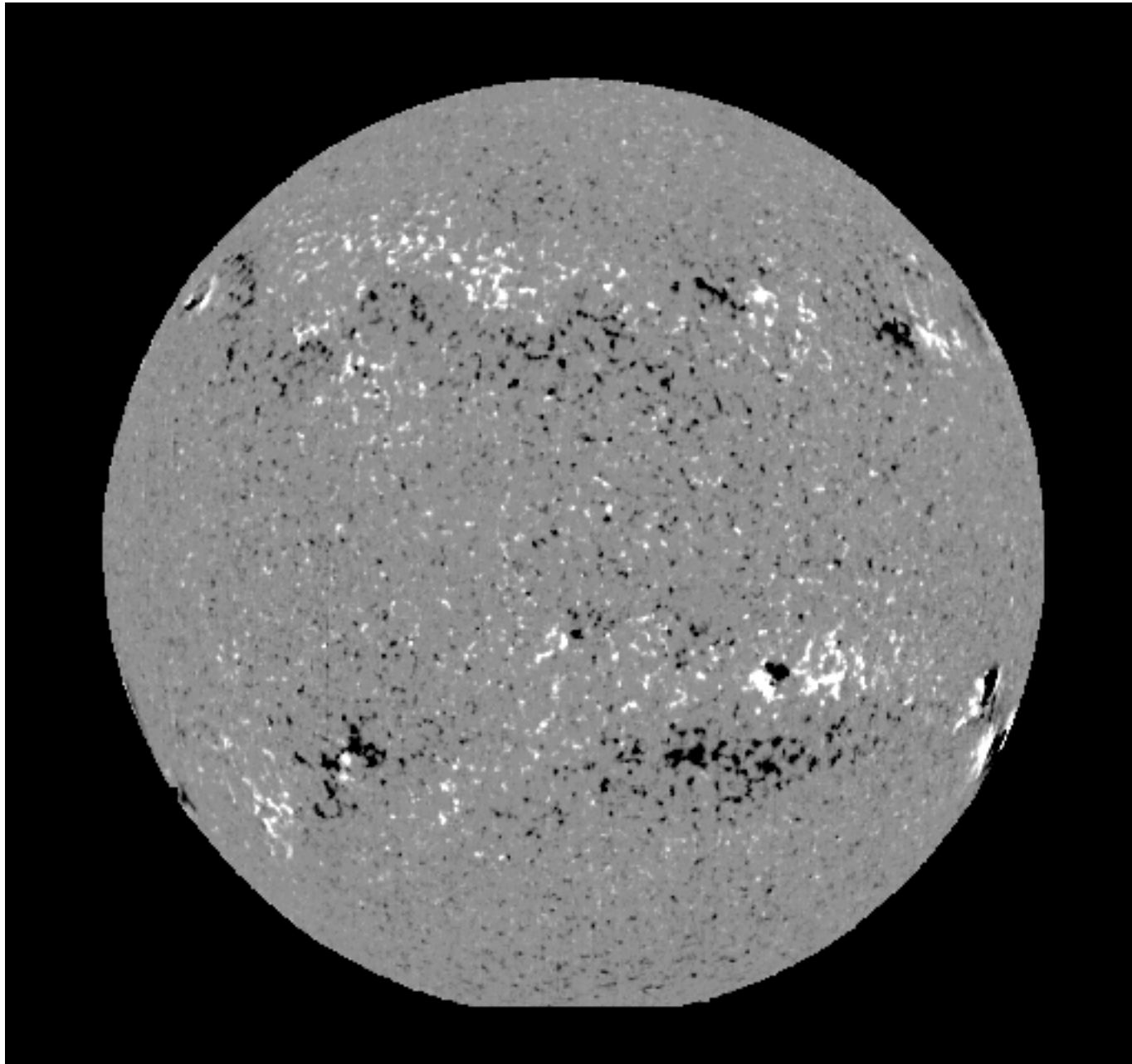
Formation of Rossby wave and Alfvén waves in the coronal hole and boundary of coronal hole region



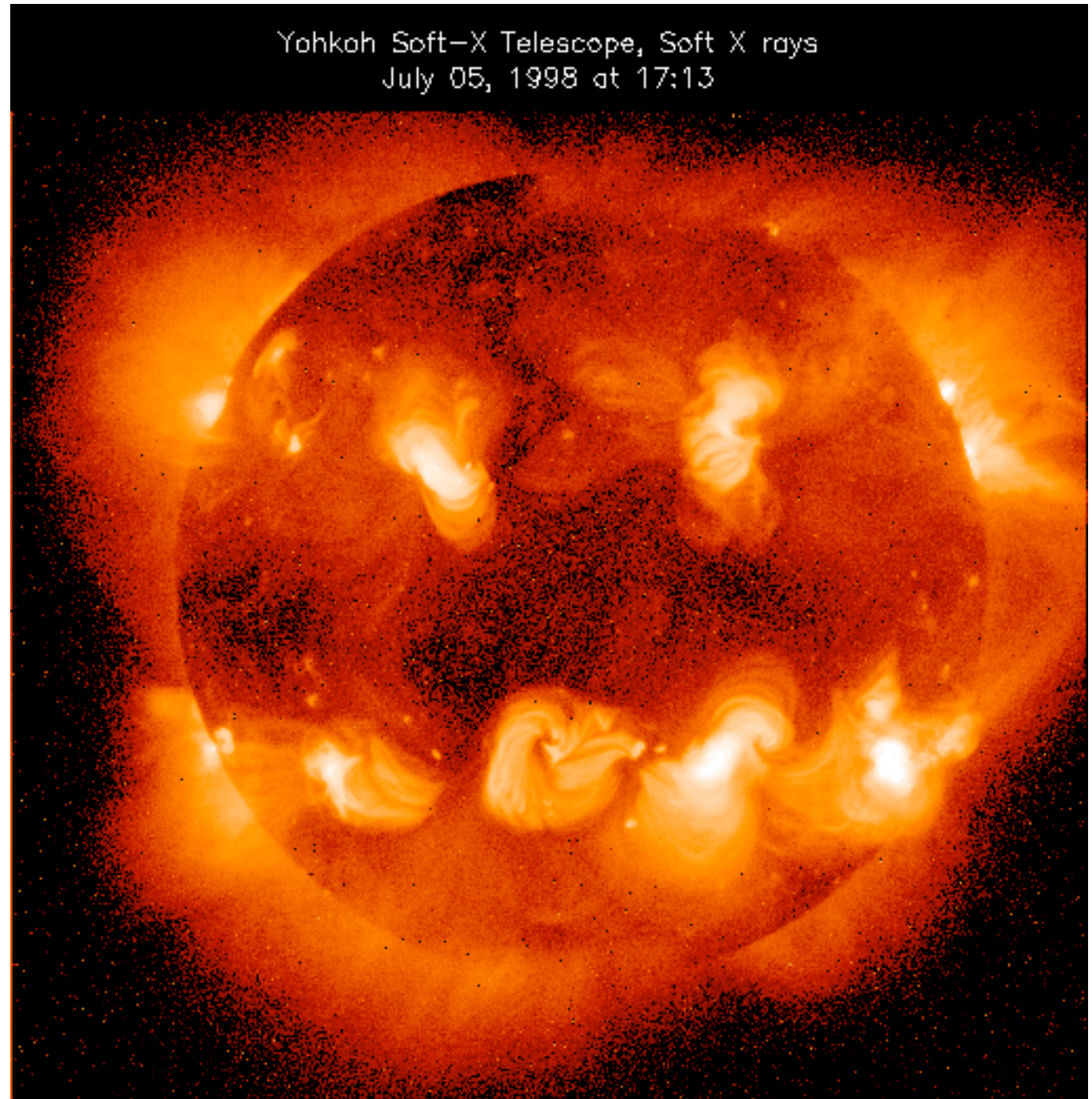
Rossby wave: $H\alpha$ observation of solar filaments (02/22/1998)



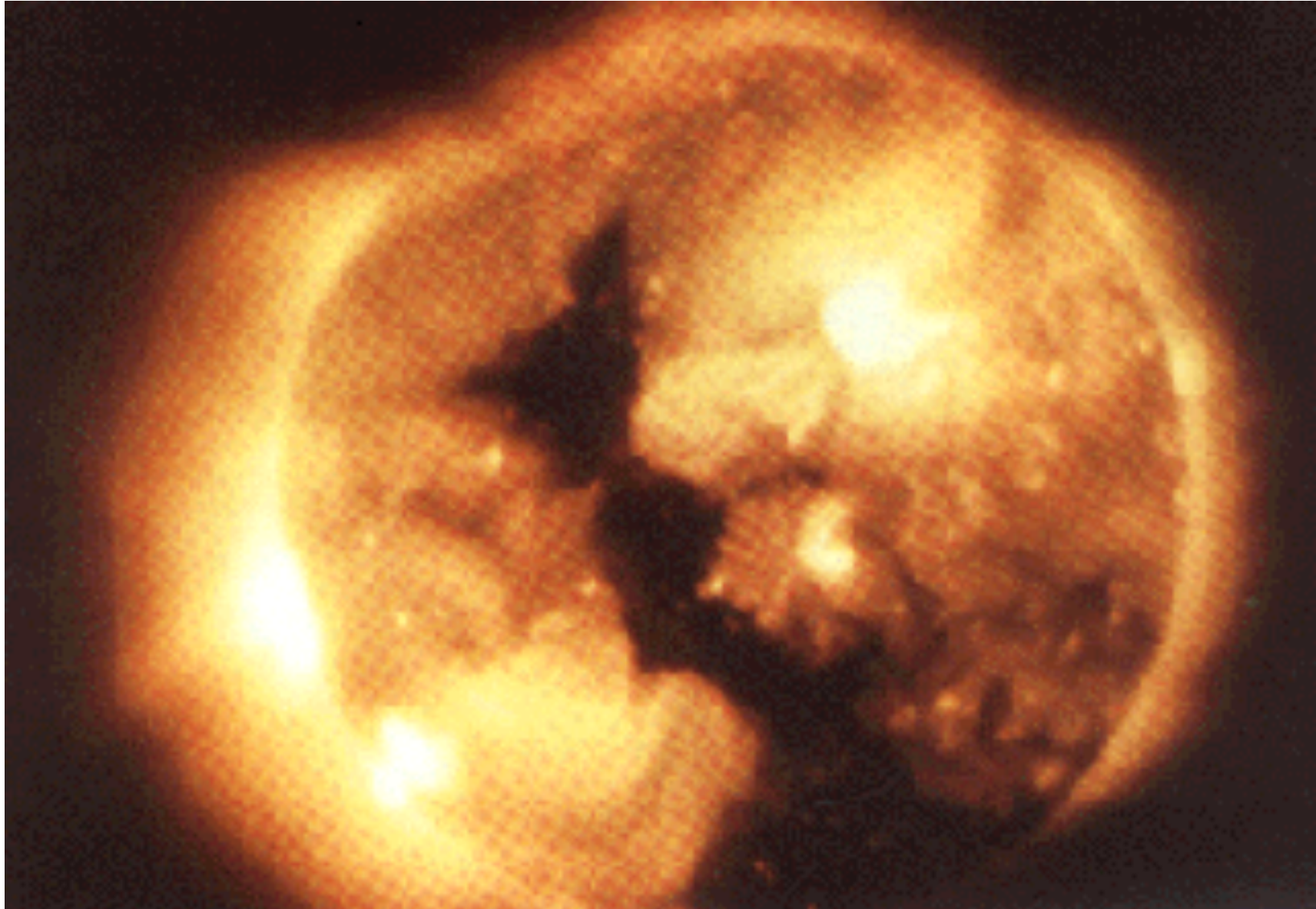
Rossby wave: MDI observation of solar magnetic field (02/22/1998)

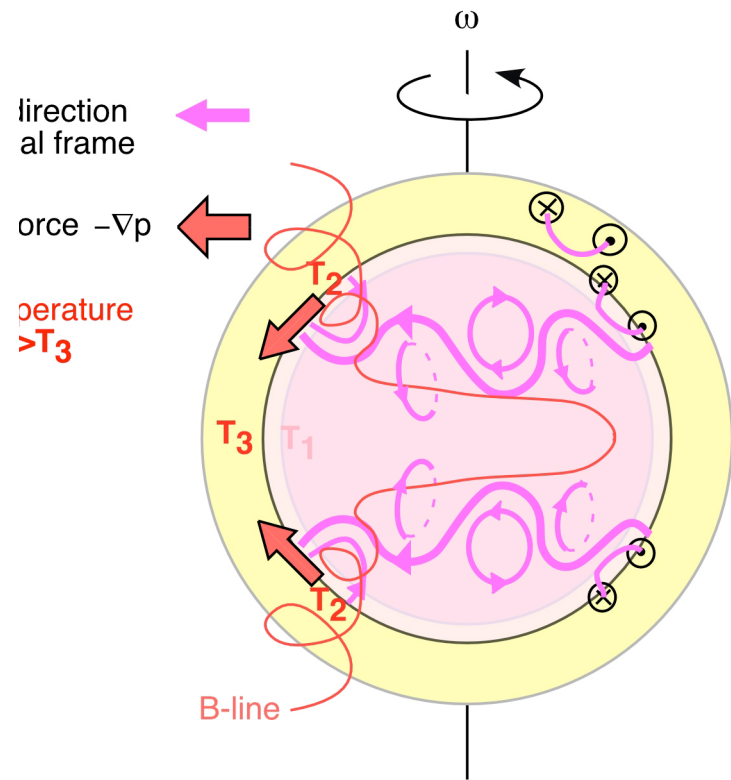


**Rossby
wave:
Yohkoh SXT
observation
07/05/1998**



Nonlinear Evolution of Rossby Wave Lead to Formation of Elephant-trunk Coronal Hole

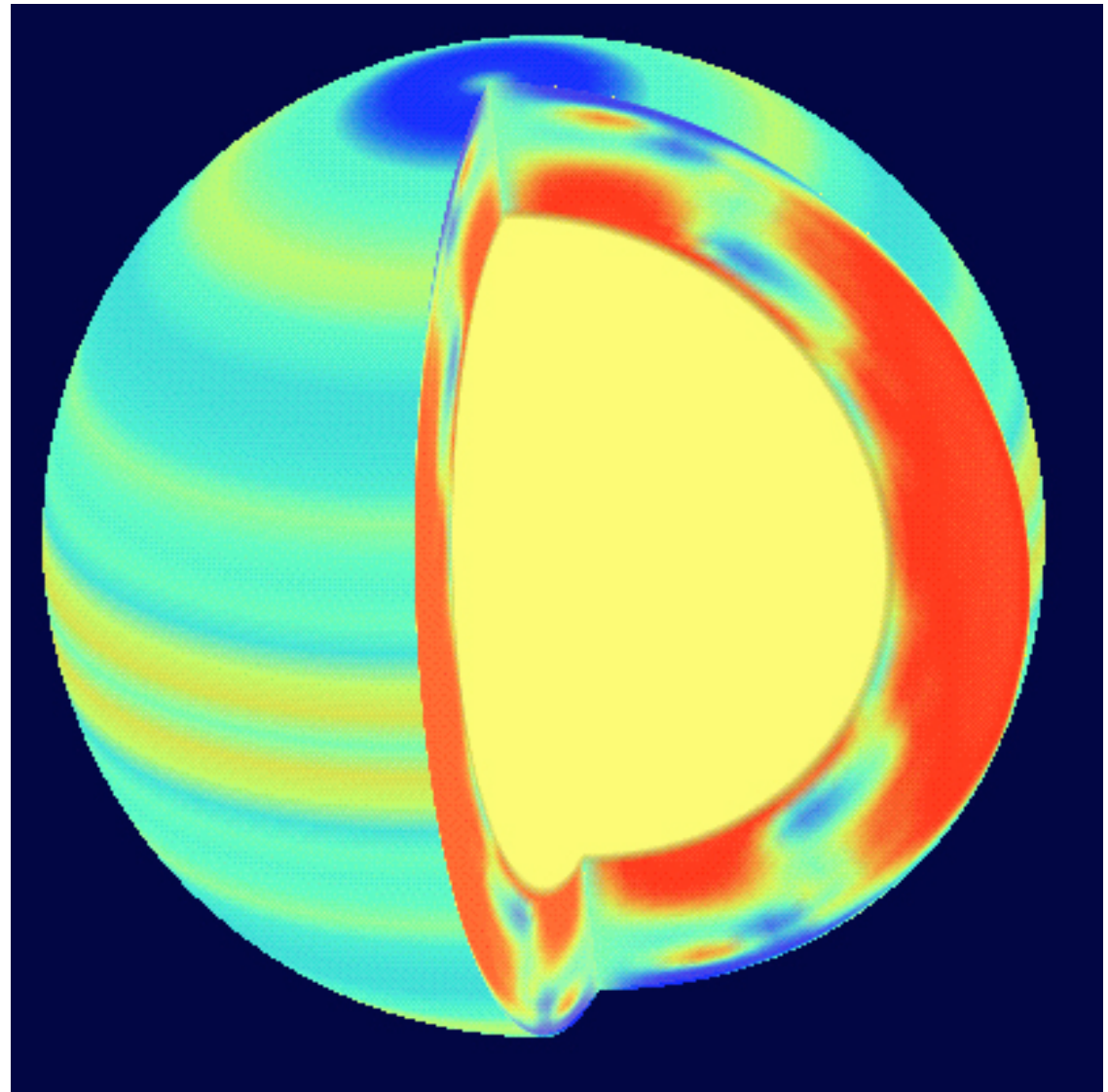




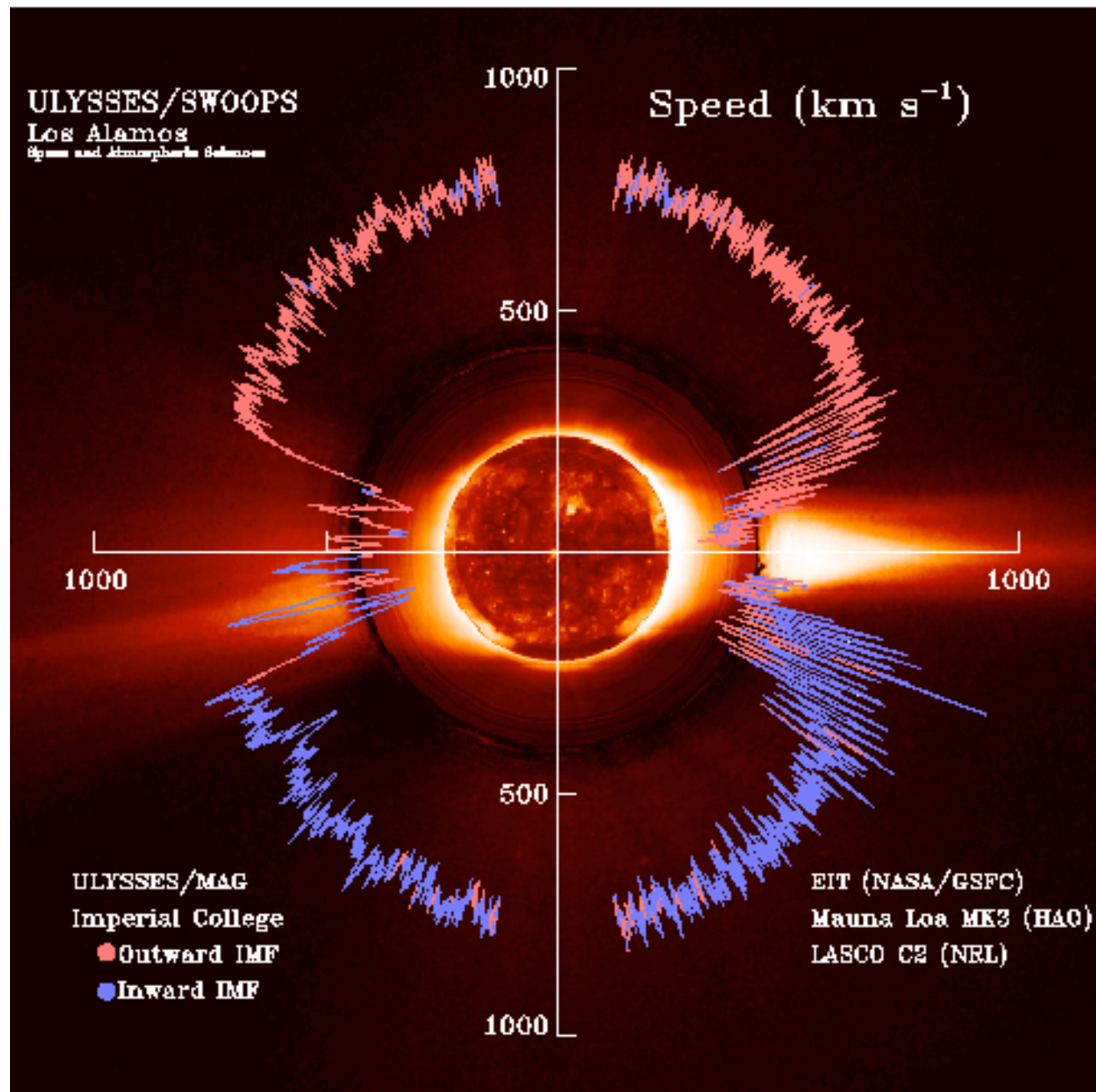
Generation of large amplitude, nearly circularly polarized Alfvén waves in the coronal hole region

MDI Observation of Solar Differential Rotation

Red: faster, Blue: slower



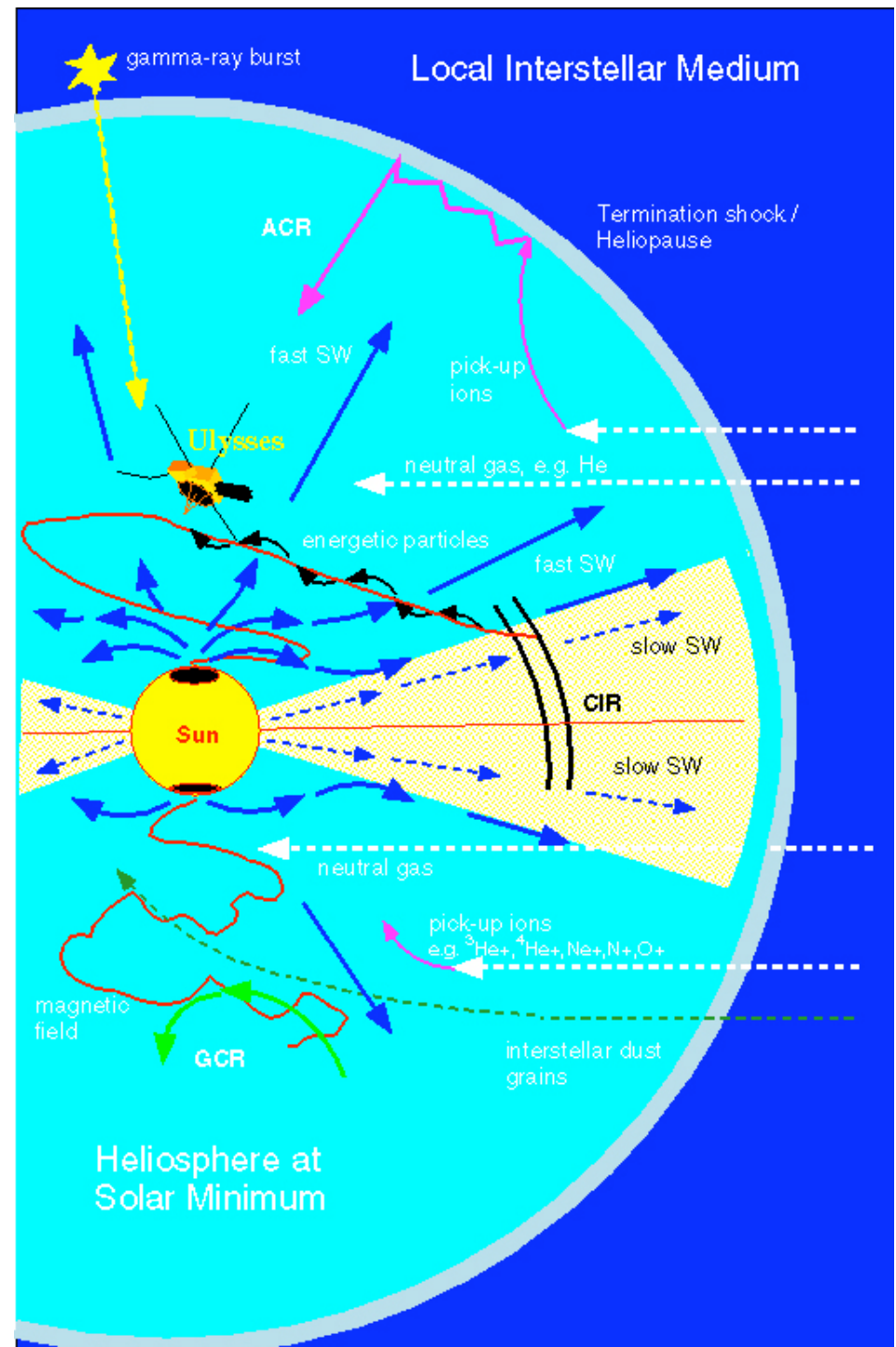
Ulysses observation of solar wind speed(1991-1998)



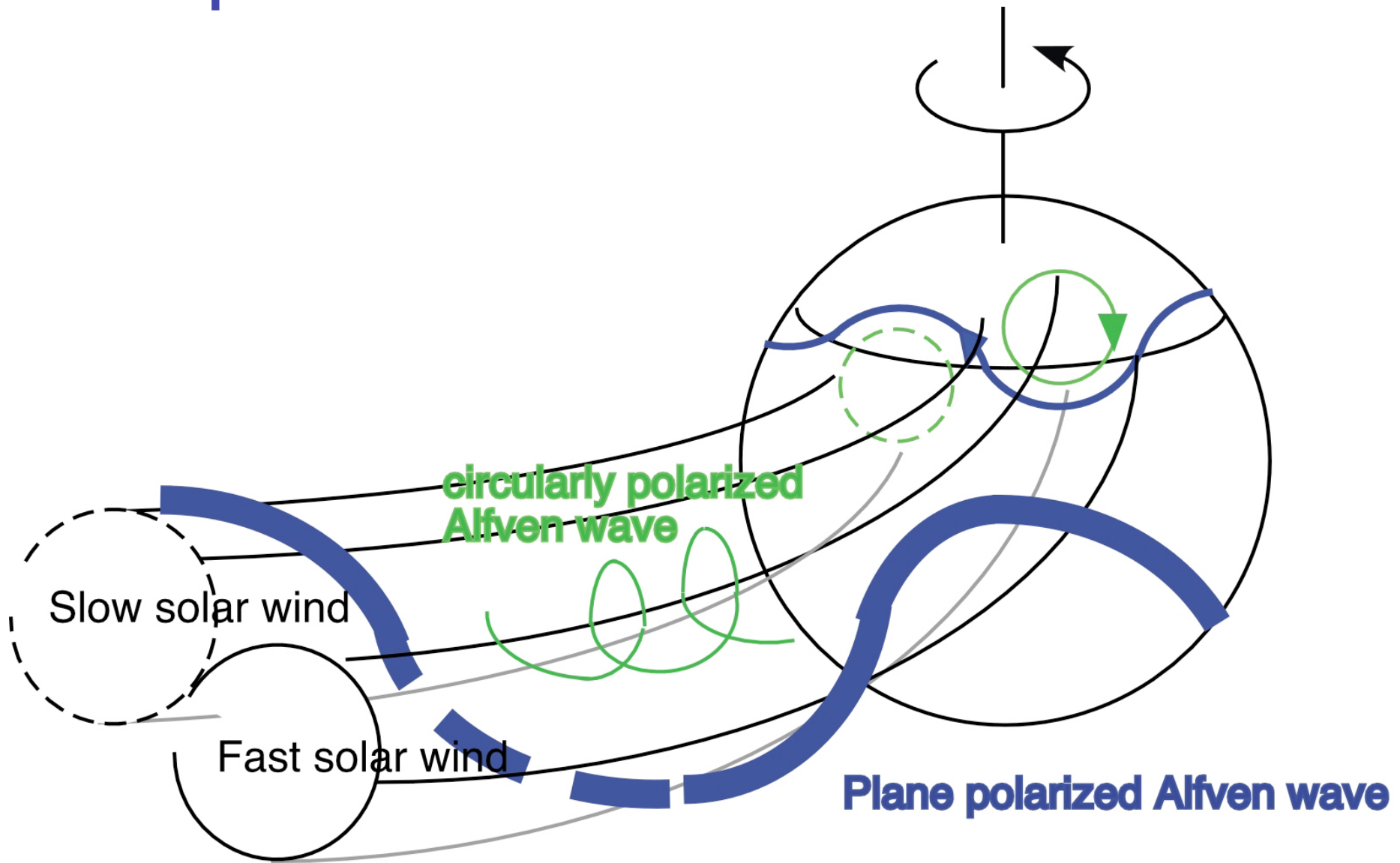
It is believed that ...
observed large-amplitude
Alfven waves are
responsible for solar wind
acceleration in the polar
coronal hole

Q: What is the energy source
of these large-amplitude
Alfven waves?

A: Rossby wave!



Rossby wave on the Sun can result in plane polarized Alfven wave in the solar wind



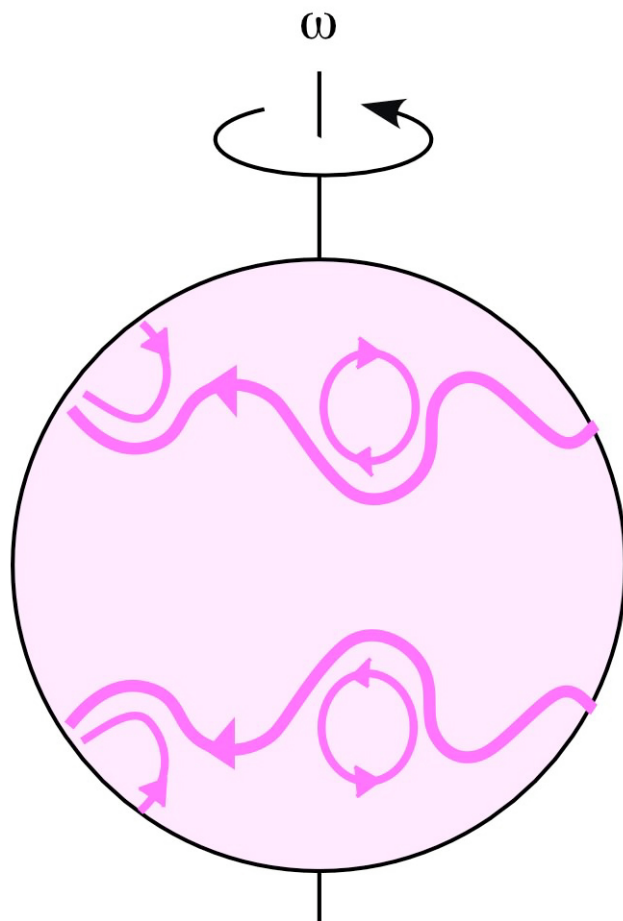
Theoretical Solutions of Steady State Nonlinear Alfven Waves and Rotational Structure in Ion-Electron Two-Fluid Plasma

(Lyu and Kan, 1989)

(Lyu, 1991)

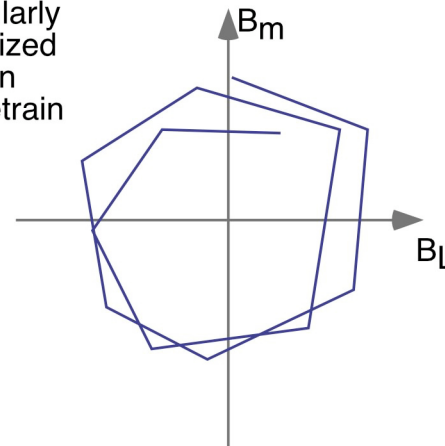
Observation of Plane Polarized Alfven Waves in the Solar Wind

--- Evidence of Rossby wave on the Sun

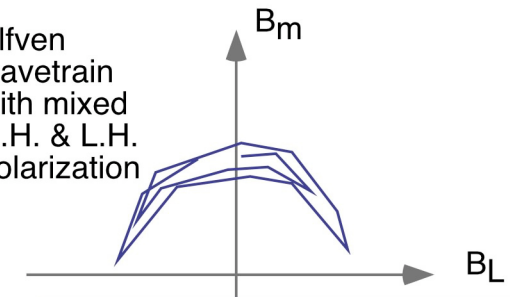


Magnetic hodograms of Alfven waves and rotational discontinuities (RDs)

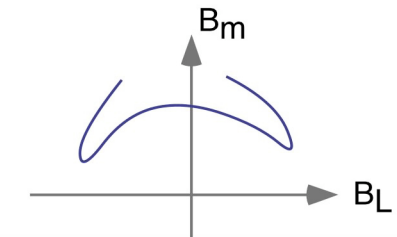
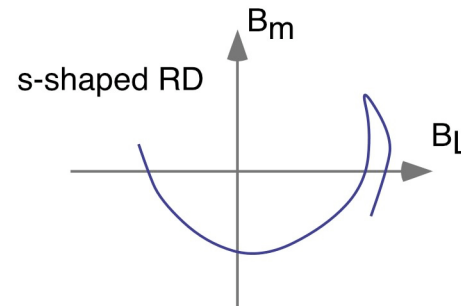
Circularly polarized Alfven wavetrain



Alfven wavetrain with mixed R.H. & L.H. polarization

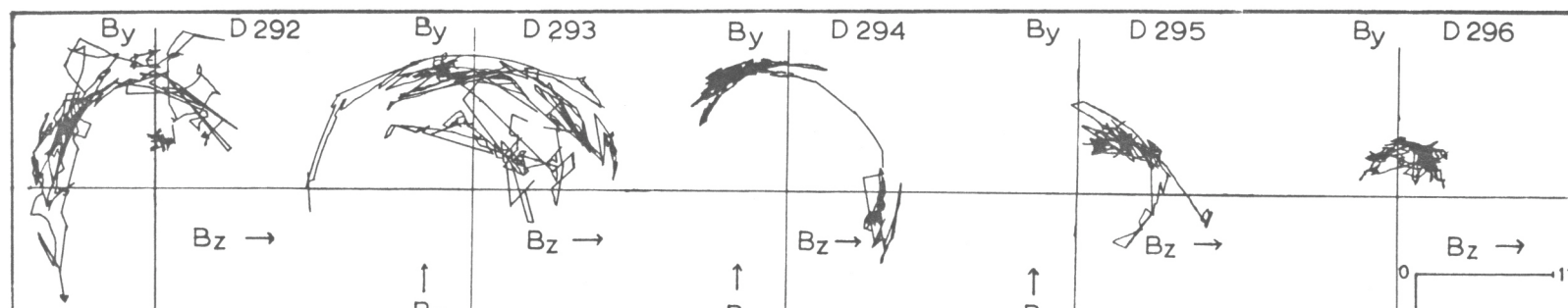
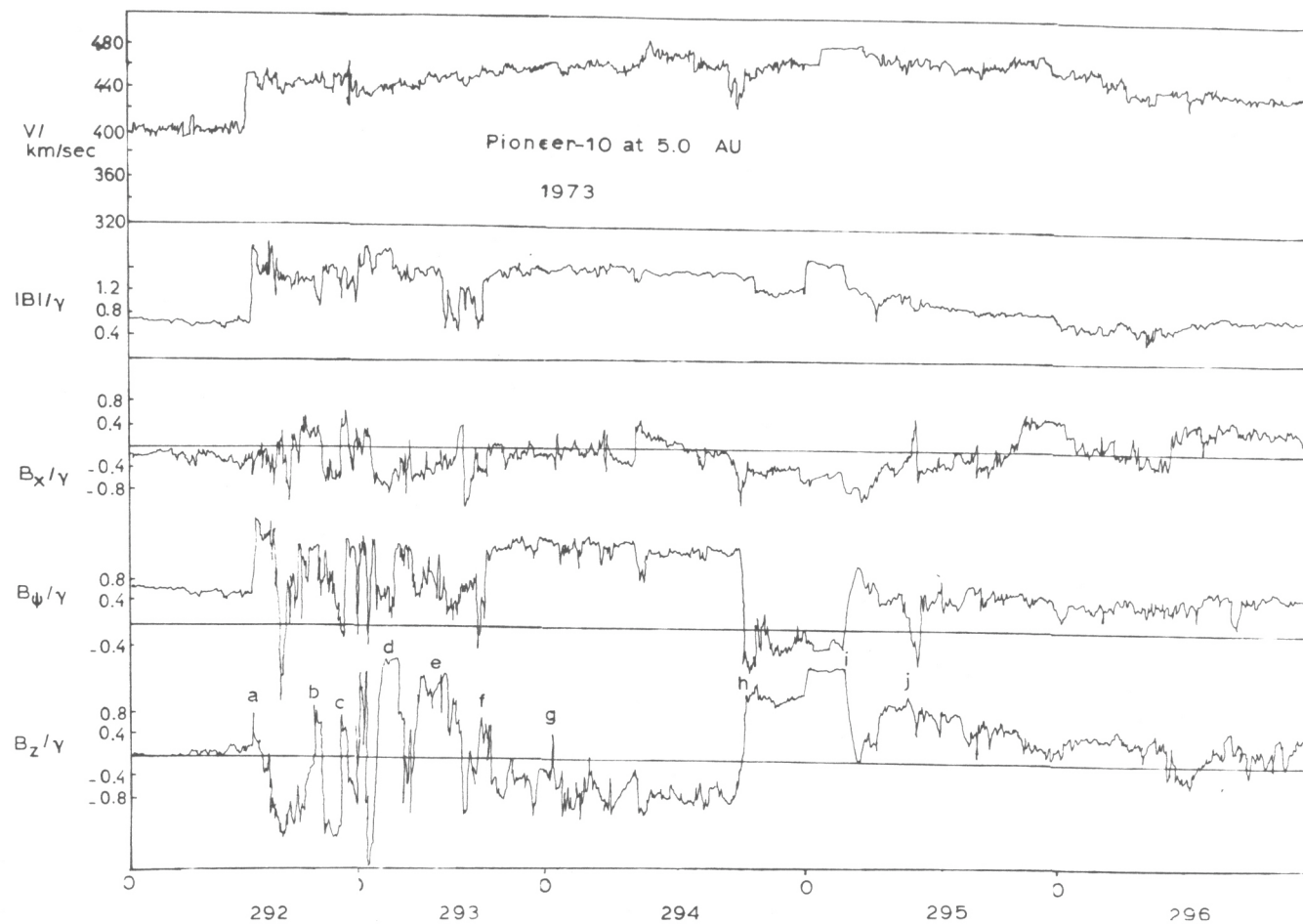


Planar polarized Alfven waves (i.e., waves with mixed polarization) have been observed by Pioneer 10 and Pioneer 11 at 5 AU and 2.5 AU, respectively. (Mavromichalaki et al., 1988)

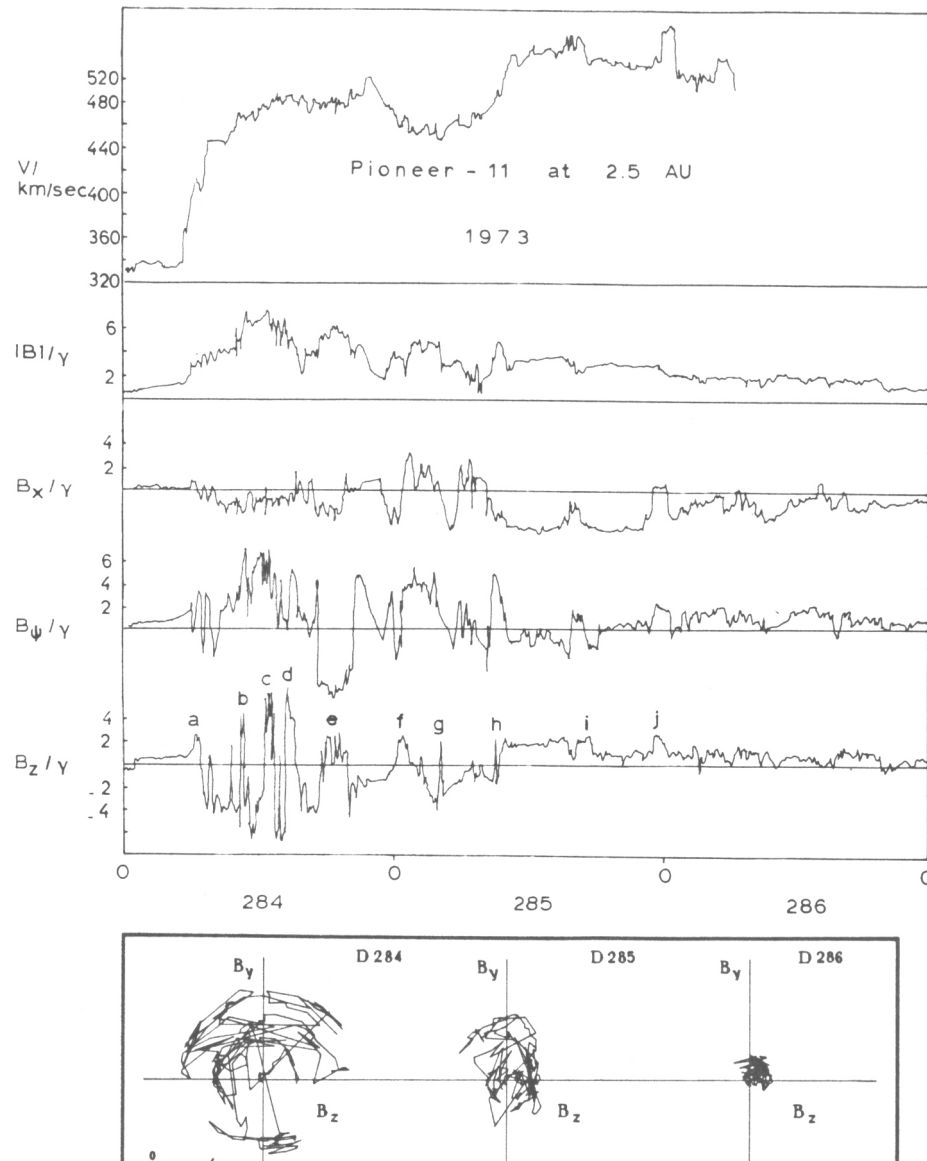


double-hooked RD (e.g., Neugebauer and Buti, 1990)

(1988)

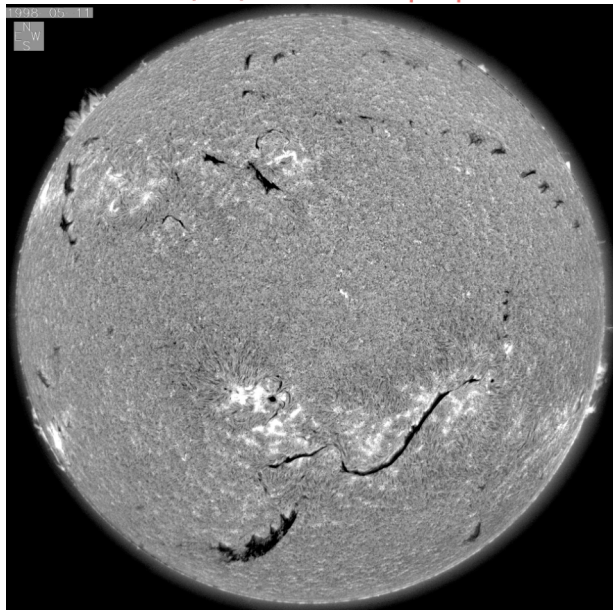
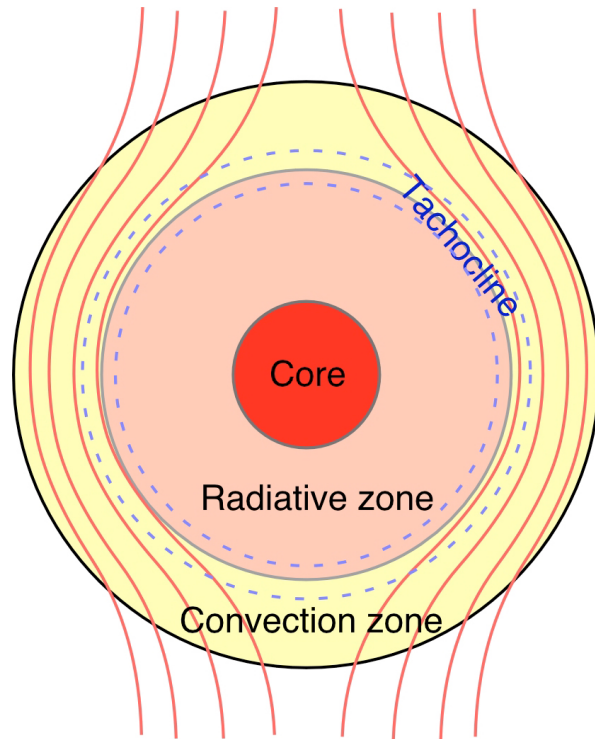


RELATIVELY STABLE, LARGE-AMPLITUDE WAVES SEEN AT 2.5 AND 5.0 AU

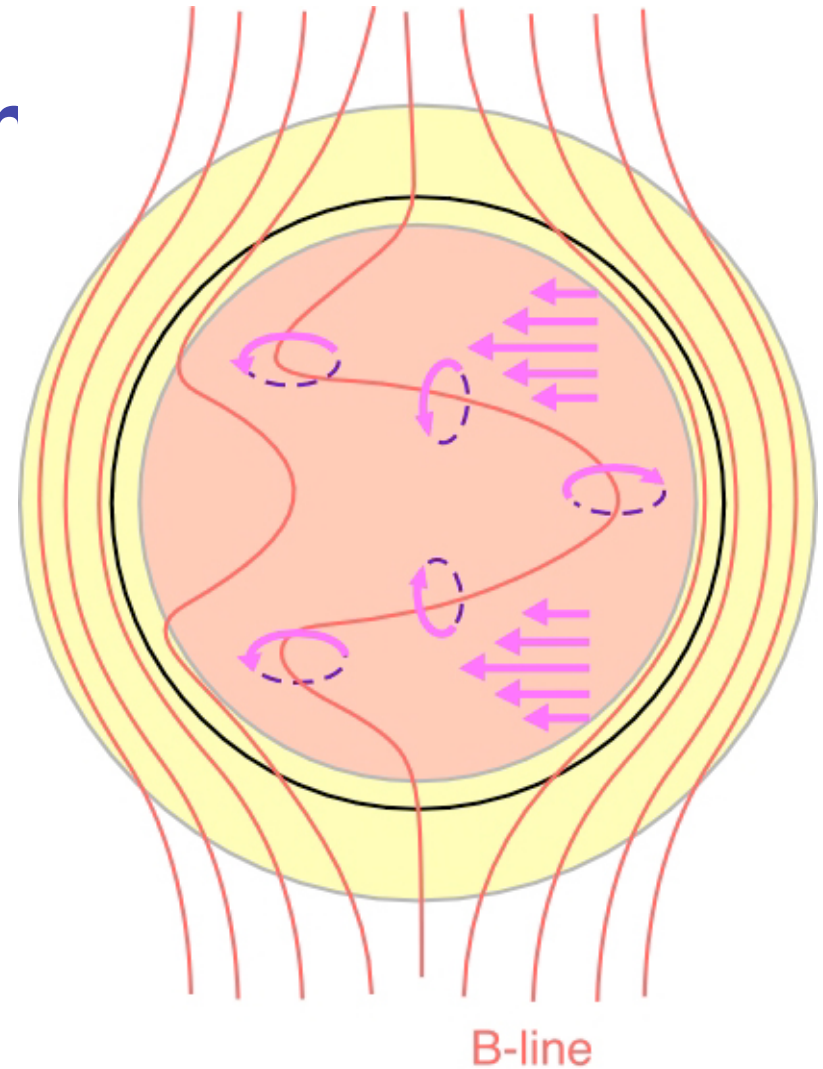


Mavromichalaki et al., 1988

Conclusior



BBO H α observation 05/11/1998

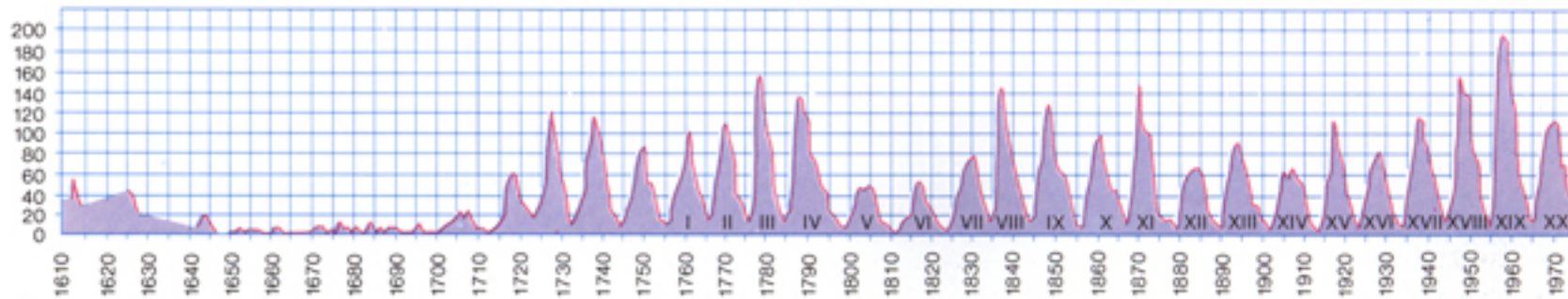


Our Model:

Plasma in solar convection zone has
Higher Temperature in the Polar Region
Easterly Wind ==> Differential Rotation

Discussion:

Cause of Maunder Minimum and variations in solar activities:



Gravitational disturbances and/or tidal effects

- >> reduce density gradient (in the tachocline region)
- >> reduce diamagnetic current
- >> reduce temperature gradient
- >> reduce differential rotation
- >> reduce solar activities
- >> reduce sunspots number

Observations: (~ 90 years Gleissberg cycle) (tidal effect)

Sunspots number decreases with increasing solar diameter.

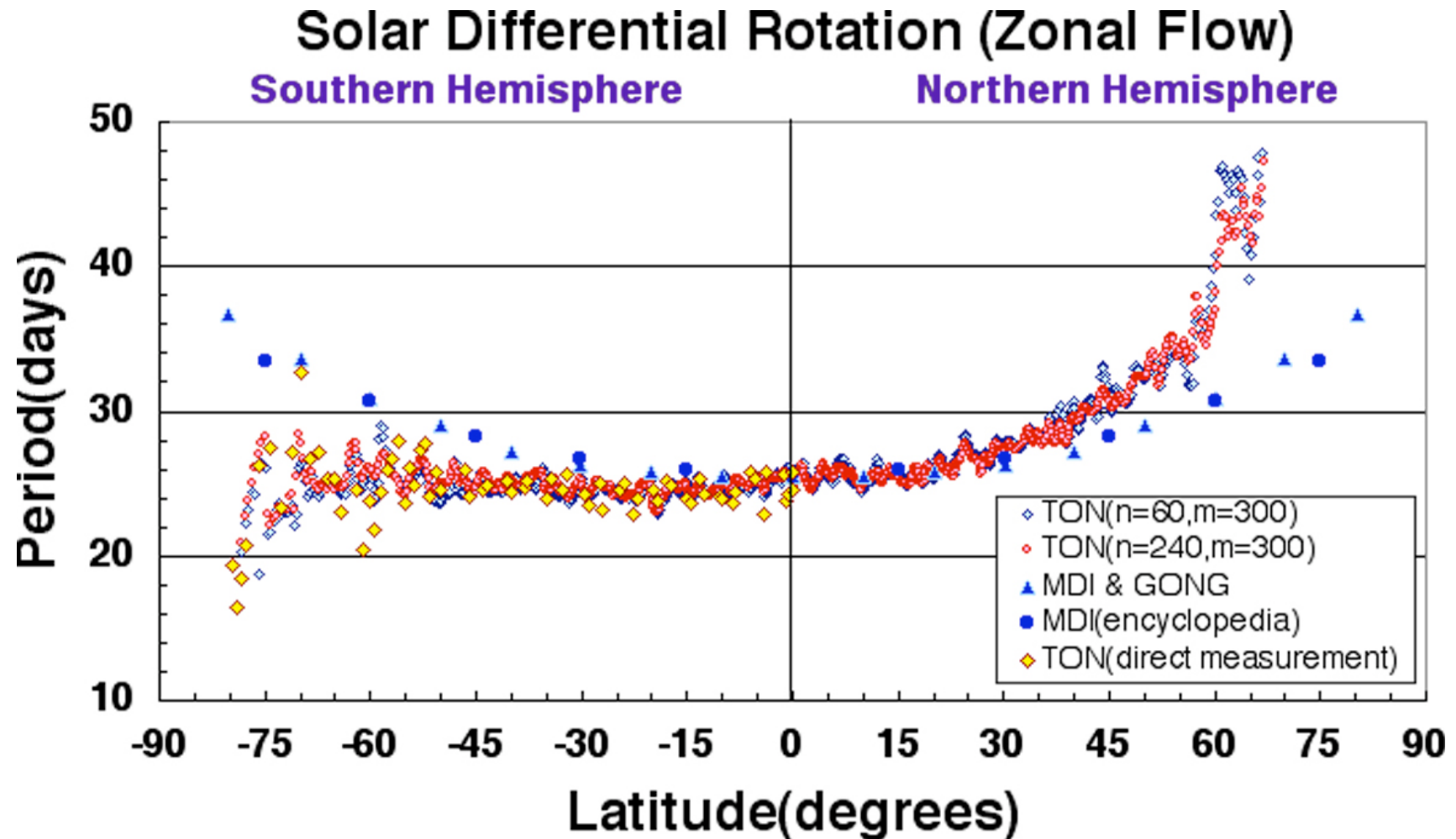
(Jupiter: 11.86 yr.) (Saturn: 29.5 yr.) (Uranus: 84 yr.)

Part 2

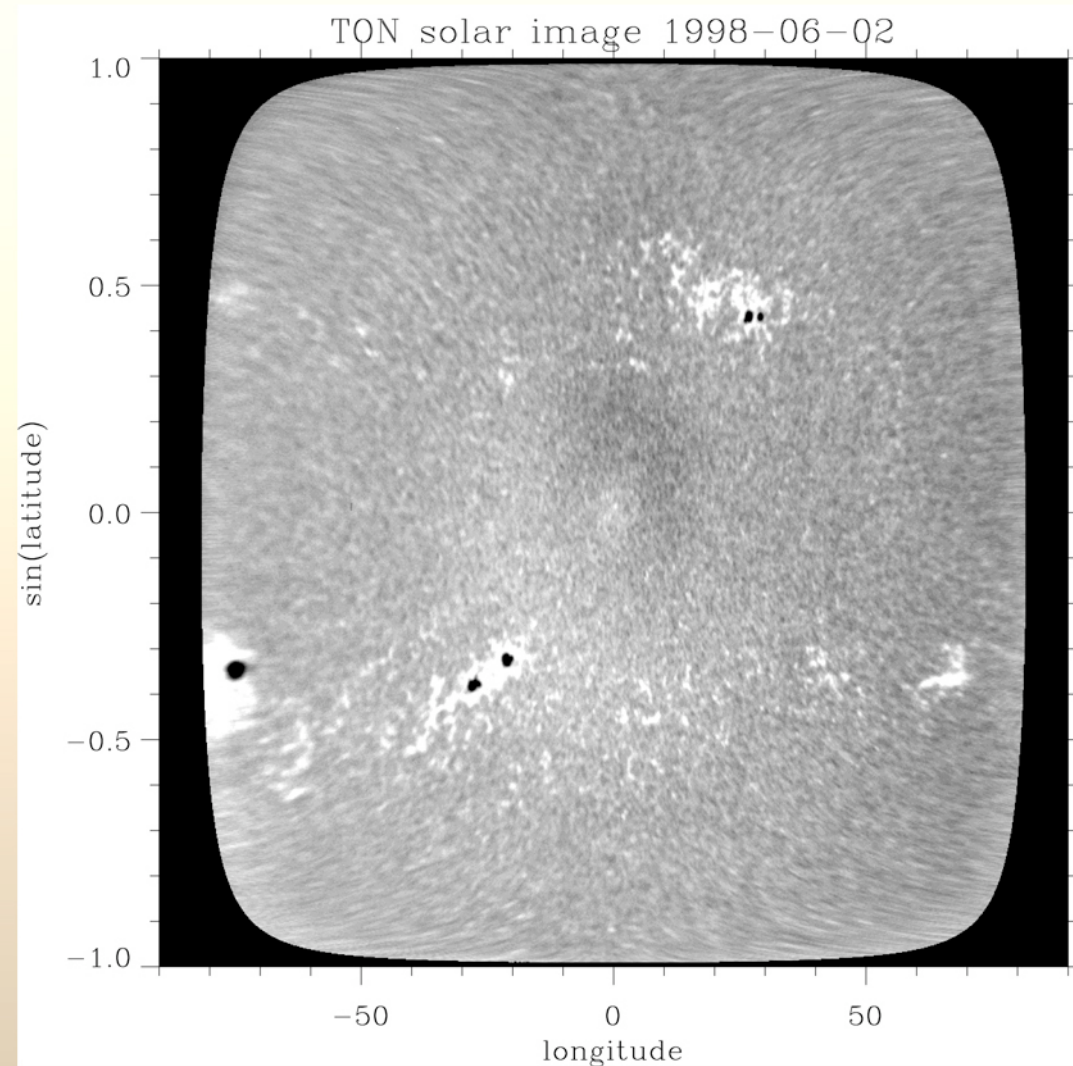
Observation of Solar Surface Differential Rotation

TON data analysis

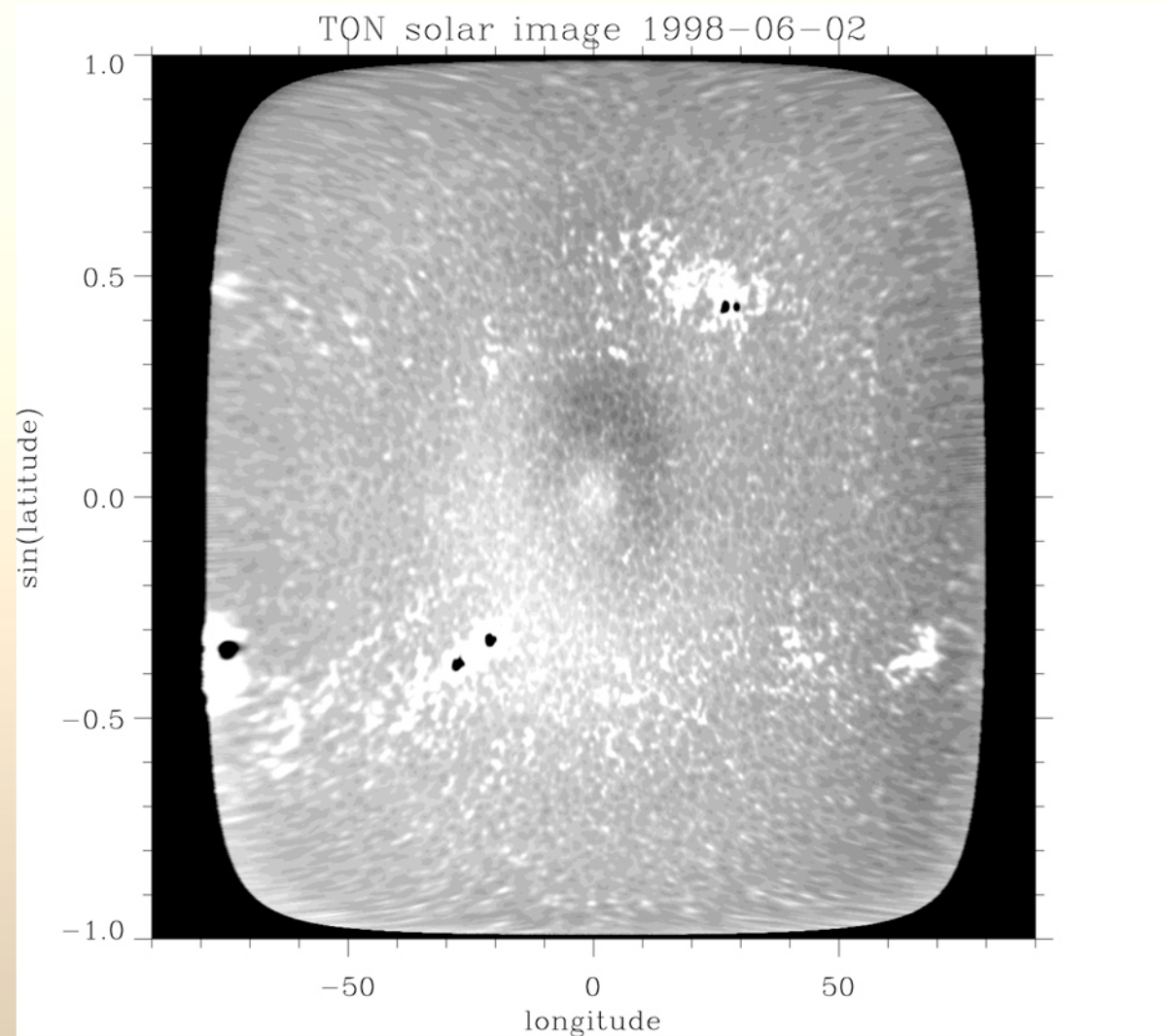
H. S. Yu, L. H. Lyu, D. Y. Chou



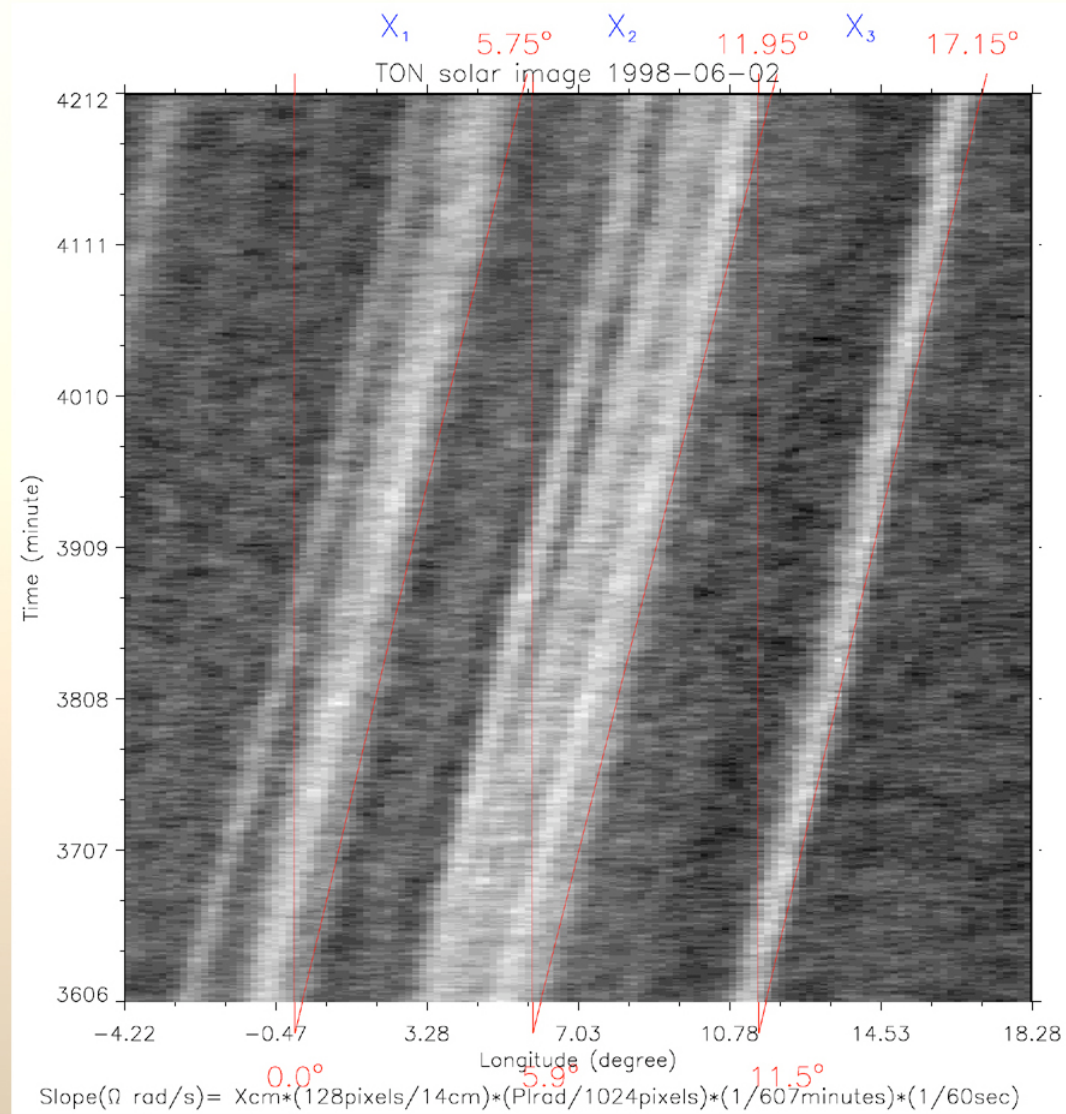
TON one-minute snapshot image (Granulations dominate)



Time integrated (>500min) image (Supergranulations dominate)



Direct measurement at latitude=27°S



Correlation analysis

Let Fourier transform of function $h(x)$ to be $h(k) = \mathcal{F}[h(x)]$,

where $h(x) \in R$ and $h(k) \in C$. It can be shown that

if $h(x) = \int f(x')g(x-x')dx'$ then $h(k) = f(k) * g(k)$.

Likewise, if $h(x) = \int f(x')g(x'+x)dx'$ then

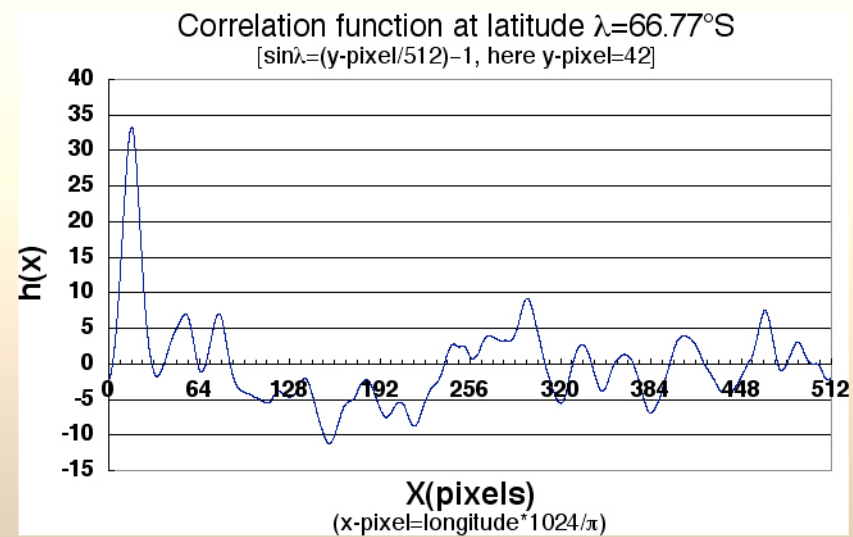
$$h^c(k) = f(k) * g^c(k) \quad (1)$$

where $h^c(k)$ is the complex conjugate of $h(k)$.

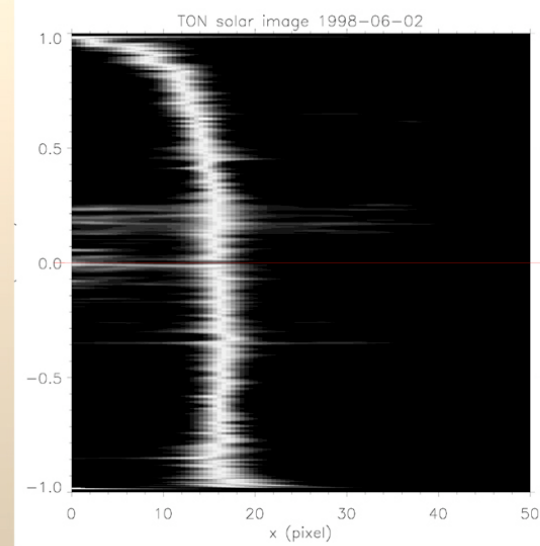
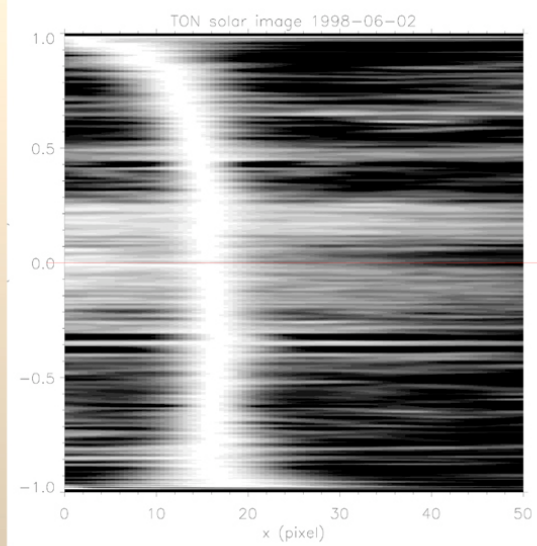
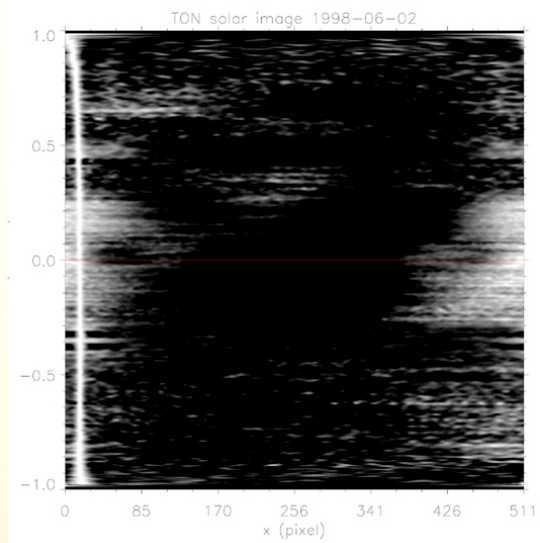
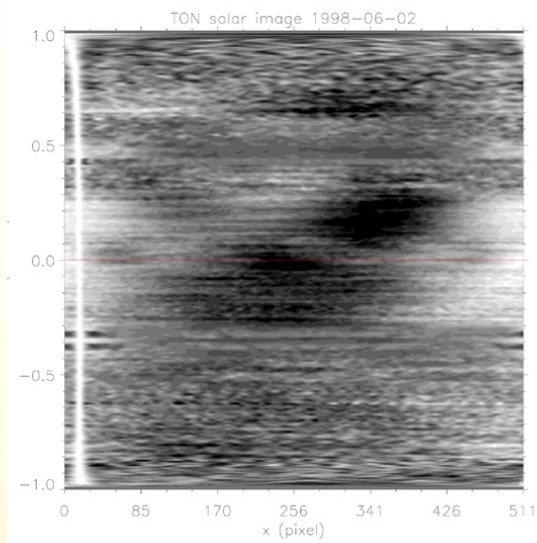
To reduce numerical error, in our correlation analysis, data of given latitude and time at different longitude are shift and normalized based on their mean and standard deviation, respectively. Autocorrelation function $h(x)$ of given latitude can be obtained from $h^c(k)$, which can be obtained from Eq. (1).

We normalized correlation function $h(x)$ at different latitude by the maximum value $h(x)$ at the given latitude. We can then plot an image of the normalized $h(x)$ at different latitude. The peak distribution of $h(x)$ can be seen from the follow four panels.

$h(x)$ before normalization



Normalized $h(x)$



Solar Differential Rotation (Zonal Flow)

