

## Appendix A. Curvature Drift

Consider a charge particle with constant magnetic moment and non-zero velocity component parallel to the local magnetic field. If curvature of the magnetic field line is non-zero, then the particle's field-aligned moving frame will become a non-inertial frame. If the parallel speed  $v_{\parallel}$  is nearly a constant, the *equation of motion* in the  $v_{\parallel}$  non-inertial moving frame, can be approximately written as

$$m \frac{d\mathbf{v}}{dt} = \frac{\hat{R}_B m v_{\parallel}^2}{R_B} + q\mathbf{v} \times \mathbf{B} \quad (\text{A.1})$$

Where  $\mathbf{v} = \mathbf{v}_{gyro} + \mathbf{v}_{drift}$ ,  $\mathbf{v}_{gyro}$  is the high frequency gyro motion velocity, and  $\mathbf{v}_{drift}$  is the low frequency (or time independent) drift velocity. Averaging Eq. (A.1) over one gyro period ( $\tau = 2\pi / \Omega_c$ , where  $\Omega_c = |q|B/m$ ), we can obtain equation for low frequency guiding-center motion in the  $v_{\parallel}$  non-inertial moving frame

$$\frac{\hat{R}_B m v_{\parallel}^2}{R_B} + q\mathbf{v}_{drift} \times \mathbf{B} = 0 \quad (\text{A.2})$$

where  $\mathbf{v}_{drift}$  is the curvature drift velocity. The curvature drift velocity can be obtained from Eq. (A.2),

$$\mathbf{v}_{drift} = \frac{\frac{\hat{R}_B m v_{\parallel}^2}{R_B} \times \mathbf{B}}{q R_B} = \frac{m v_{\parallel}^2}{q B^2} (\hat{R}_B \times \mathbf{B}) \quad (\text{A.3})$$

Since

$$\begin{aligned} -\frac{\hat{R}_B}{R_B} &= \frac{d\hat{B}}{ds} = \hat{B} \cdot \nabla \hat{B} = \frac{\mathbf{B}}{B} \cdot \nabla \left( \frac{\mathbf{B}}{B} \right) = \frac{\mathbf{B} \cdot \nabla \mathbf{B}}{B^2} - \frac{\mathbf{B} \mathbf{B} \cdot \nabla B}{B^3} \\ &= \frac{-\mathbf{B} \times (\nabla \times \mathbf{B}) + \nabla \frac{B^2}{2}}{B^2} - \frac{\hat{B} \hat{B} \cdot \nabla B}{B} = \frac{-\mathbf{B} \times (\nabla \times \mathbf{B})}{B^2} + \frac{\nabla B}{B} - \frac{\hat{B} \hat{B} \cdot \nabla B}{B} \\ &= \frac{-\mathbf{B} \times (\nabla \times \mathbf{B})}{B^2} + (\mathbf{1} - \hat{B} \hat{B}) \cdot \frac{\nabla B}{B} \end{aligned} \quad (\text{A.4})$$

where

$$\mathbf{1} = \hat{x}\hat{x} + \hat{y}\hat{y} + \hat{z}\hat{z} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

is the unit tensor. For any given vector  $\mathbf{A}$ , we have  $\mathbf{A} = \mathbf{A} \cdot \mathbf{1} = \mathbf{1} \cdot \mathbf{A}$ .

If the magnetic field is along the z-direction, then  $\hat{B} \hat{B} = \hat{z} \hat{z}$ , and

$$\mathbf{1} - \hat{B}\hat{B} = \hat{x}\hat{x} + \hat{y}\hat{y} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Any vector  $\mathbf{A}$  can be decomposed into two parts, one is parallel to the local magnetic field ( $\mathbf{A}_{\parallel}$ ), another is perpendicular to the local magnetic field ( $\mathbf{A}_{\perp}$ ). Namely,  $\mathbf{A} = \mathbf{A}_{\parallel} + \mathbf{A}_{\perp}$ . It can be shown that

$$\mathbf{A}_{\parallel} = \hat{B}\hat{B} \cdot \mathbf{A}$$

and

$$\mathbf{A}_{\perp} = (\mathbf{1} - \hat{B}\hat{B}) \cdot \mathbf{A}$$

Likewise, we have

$$\nabla_{\perp} B = \nabla B - \hat{B}\hat{B} \cdot \nabla B = (\mathbf{1} - \hat{B}\hat{B}) \cdot \nabla B \quad (\text{A.5})$$

Substituting Eq. (A.5) into Eq. (A.4), it yields

$$-\frac{\hat{R}_B}{R_B} = \frac{-\mathbf{B} \times (\nabla \times \mathbf{B})}{B^2} + \frac{\nabla_{\perp} B}{B}$$

or

$$\frac{\hat{R}_B}{R_B} = \frac{\mathbf{B} \times (\nabla \times \mathbf{B})}{B^2} - \frac{\nabla_{\perp} B}{B} \quad (\text{A.6})$$

Substituting Eq. (A.6) into Eq. (A.3), it yields

$$\begin{aligned} \mathbf{v}_{drift} &= \frac{mv_{\parallel}^2}{qB^2} \left( \frac{\hat{R}_B}{R_B} \times \mathbf{B} \right) \\ &= \frac{mv_{\parallel}^2}{qB^2} \left[ \left( \frac{\mathbf{B} \times (\nabla \times \mathbf{B})}{B^2} - \frac{\nabla_{\perp} B}{B} \right) \times \mathbf{B} \right] \\ &= \frac{mv_{\parallel}^2}{qB^2} \left[ \frac{B^2 (\nabla \times \mathbf{B})}{B^2} - \frac{\mathbf{B}\mathbf{B} \cdot (\nabla \times \mathbf{B})}{B^2} - \frac{\nabla_{\perp} B}{B} \times \mathbf{B} \right] \\ &= \frac{mv_{\parallel}^2}{qB^2} \left[ (\mathbf{1} - \hat{B}\hat{B}) \cdot (\nabla \times \mathbf{B}) - \frac{\nabla_{\perp} B}{B} \times \mathbf{B} \right] \\ &= \frac{mv_{\parallel}^2}{qB^2} \left[ (\nabla \times \mathbf{B})_{\perp} - \frac{\nabla_{\perp} B}{B} \times \mathbf{B} \right] \end{aligned}$$

Namely, the curvature drift velocity given in (A.3) can be written as

$$\mathbf{v}_{drift} = \frac{mv_{\parallel}^2}{qB^2} \left[ (\nabla \times \mathbf{B})_{\perp} - \frac{\nabla_{\perp} B}{B} \times \mathbf{B} \right] \quad (\text{A.7})$$