

Motion of a Particle in a Two- or Three-Dimensional System

PART-A

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Symon (1960): Chapter 3: Motion of a Particle in Two or Three Dimensions

“二元”或“三元”「聯立的」微分方程式 (System ODEs)

Part 1:

3-1 Vector algebra

3-3 Differentiation and integration of vectors

3-2 Applications to a set of forces acting on a particle

3-17 Motion of a particle in an electromagnetic field

帶電粒子在靜電場中的運動

帶電粒子在磁場中的迴旋運動，

帶電粒子在靜電場與磁場中的運動 (如何求解：聯立ODEs)

(如何由運動軌跡，猜可能的答案)

Vector algebra

- 向量相加
- 向量相減
- 向量相乘
 - 向量內積
 - 向量外積
 - 向量平行積
- 張量
 - 純量 (零階張量)
 - 向量 (一階張量)
 - 二階張量
 - ...

Differentiation of Vectors

$$\frac{d\vec{A}}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\vec{A}(t + \Delta t) - \vec{A}(t)}{\Delta t}$$

$$\frac{d}{dt} (\vec{A} + \vec{B}) = \frac{d}{dt} (\vec{A}) + \frac{d}{dt} (\vec{B})$$

$$\frac{d}{dt} (f\vec{A}) = \frac{df}{dt} \vec{A} + f \frac{d}{dt} (\vec{A})$$

$$\frac{d}{dt} (\vec{A} \cdot \vec{B}) = \frac{d\vec{A}}{dt} \cdot \vec{B} + \vec{A} \cdot \frac{d\vec{B}}{dt}$$

$$\frac{d}{dt} (\vec{A} \times \vec{B}) = \frac{d\vec{A}}{dt} \times \vec{B} + \vec{A} \times \frac{d\vec{B}}{dt}$$

$$\frac{d}{dt} (\vec{A}\vec{B}) = \frac{d\vec{A}}{dt} \vec{B} + \vec{A} \frac{d\vec{B}}{dt}$$

Integration of Vectors

$$W = \int \vec{F} \cdot d\vec{r} = \int \vec{F} \cdot \left(\frac{d\vec{r}}{ds} \right) ds$$

For

$$\begin{aligned} \vec{F} &= \hat{x}F_x + \hat{y}F_y + \hat{z}F_z \\ \frac{d\vec{r}}{ds} &= \hat{x} \frac{dx}{ds} + \hat{y} \frac{dy}{ds} + \hat{z} \frac{dz}{ds} \end{aligned}$$

It yields, for a given trajectory $[x(s), y(s), z(s)]$

$$W = \int \vec{F} \cdot d\vec{r} = \int \left[F_x \frac{dx}{ds} + F_y \frac{dy}{ds} + F_z \frac{dz}{ds} \right] ds$$

Integration of Vectors

$$W = \int \vec{F} \cdot d\vec{r} = \int \vec{F} \cdot \left(\frac{d\vec{r}}{dt} \right) dt = \int (\vec{F} \cdot \vec{v}) dt$$

Thus

$$\frac{dW}{dt} = \vec{F} \cdot \vec{v} = \text{Power}$$

Non-relativistic Particle Motion

$$\frac{d\vec{x}}{dt} = \vec{v}$$

$$\frac{d\vec{v}}{dt} = \vec{a}$$

$$\vec{F} = m\vec{a} = m\frac{d\vec{v}}{dt}$$

$$\vec{p} = m\vec{v}$$

$$T = \frac{1}{2}mv^2 = \frac{1}{2}m\vec{v} \cdot \vec{v}$$

Motion of a Charged Particle

- Let us consider a charge particle with mass m and charge q .
- Let us assume that the Lorentz force is much stronger than the the gravitational force, we shall ignore the gravitational force in the following discussion.

- 帶電粒子在均勻的靜電場 $\vec{E} = \hat{y}E_y$ 中的運動

$$\frac{d\vec{v}}{dt} = \frac{q}{m} \vec{E} \Rightarrow \dot{v}_y = \left(\frac{q}{m}\right) E_y$$

- 帶電粒子在均勻的磁場 $\vec{B} = \hat{z}B_z$ 中的迴旋運動

$$\frac{d\vec{v}}{dt} = \frac{q}{m} \vec{v} \times \vec{B} \Rightarrow \begin{cases} \dot{v}_x = \left(\frac{q}{m}\right) v_y B_z \\ \dot{v}_y = -\left(\frac{q}{m}\right) v_x B_z \end{cases}$$

- 帶電粒子在均勻的靜電場 $\vec{E} = \hat{y}E_y$ 與均勻的磁場 $\vec{B} = \hat{z}B_z$ 中的運動

$$\frac{d\vec{v}}{dt} = \frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \Rightarrow \begin{cases} \dot{v}_x = \left(\frac{q}{m}\right) v_y B_z \\ \dot{v}_y = \left(\frac{q}{m}\right) [E_y - v_x B_z] \end{cases}$$

Motion of a Charged Particle

- Let us consider a charge particle with mass m and charge q .
- Let us assume that the Lorentz force is much stronger than the the gravitational force, we shall ignore the gravitational force in the following discussion.
- 帶電粒子在均勻的靜電場 \vec{E} 與均勻的磁場 \vec{B} 中的運動 & $E/B \ll c$

$$\begin{aligned} \frac{d\vec{v}}{dt} &= \frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \Rightarrow \frac{d^2\vec{v}}{dt^2} = \frac{q}{m} \left(\frac{d\vec{v}}{dt} \times \vec{B} \right) \\ \Rightarrow \frac{d^2\vec{v}}{dt^2} &= \frac{q}{m} \left[\frac{q}{m} (\vec{E} + \vec{v} \times \vec{B}) \times \vec{B} \right] = \frac{q^2}{m^2} \vec{E} \times \vec{B} + \frac{q^2}{m^2} (\vec{v} \times \vec{B}) \times \vec{B} \\ \Rightarrow \frac{d^2\vec{v}}{dt^2} &= \frac{q^2 B^2}{m^2} \left[\frac{\vec{E} \times \vec{B}}{B^2} + \vec{v} \cdot \hat{B} \hat{B} - \vec{v} \right] = -\frac{q^2 B^2}{m^2} \left[\vec{v}_\perp - \frac{\vec{E} \times \vec{B}}{B^2} \right] \end{aligned}$$

where $\vec{v}_\perp = \vec{v} - \vec{v}_\parallel = \vec{v} - \vec{v} \cdot \hat{B} \hat{B} = \vec{v} - \vec{v} \cdot \frac{\vec{B}\vec{B}}{B^2}$. Let $\vec{V}_\perp = \vec{v}_\perp - \frac{\vec{E} \times \vec{B}}{B^2}$, it yields

$$\begin{aligned} \frac{d^2\vec{V}_\perp}{dt^2} &= -\frac{q^2 B^2}{m^2} \vec{V}_\perp \\ \frac{d^2\vec{v}_\parallel}{dt^2} &= 0 \end{aligned}$$

Thus, we have $\vec{v}_\perp = \frac{\vec{E} \times \vec{B}}{B^2} + \vec{V}_\perp = \vec{v}_{E \times B} + \vec{v}_{gyro}$

$$(\vec{A} \times \vec{B}) \times \vec{C} = (\vec{A} \cdot \vec{C}) \vec{B} - (\vec{B} \cdot \vec{C}) \vec{A}$$

Symon (1960): Chapter 3: Motion of a Particle in Two or Three Dimensions

“二元”或“三元”「聯立的」微分方程式 (System ODEs)

Part 2:

3-4 Kinematics in a plane

3-5 Kinematics in three dimensions

3-6 Elements of vector analysis

Kinematic Equations in Different Coordinate Systems

	Cartesian	Cylindrical	Spherical
Coord.	$[x(t), y(t), z(t)]$	$[r(t), \theta(t), z(t)]$	$[r(t), \theta(t), \phi(t)]$
position	$\vec{r}(t) = x(t)\hat{x} + y(t)\hat{y} + z(t)\hat{z}$	$\vec{r}(t) = r(t)\hat{r}(\theta(t)) + z(t)\hat{z}$	$\vec{r}(t) = r(t)\hat{r}(\theta(t), \phi(t))$
$\vec{v} = \dot{\vec{r}}$ $\vec{a} = \ddot{\vec{r}}$	$\vec{v} = v_x\hat{x} + v_y\hat{y} + v_z\hat{z}$ $\vec{a} = a_x\hat{x} + a_y\hat{y} + a_z\hat{z}$	$\vec{v} = v_r\hat{r} + v_\theta\hat{\theta} + v_z\hat{z}$ $\vec{a} = a_r\hat{r} + a_\theta\hat{\theta} + a_z\hat{z}$	$\vec{v} = v_r\hat{r} + v_\theta\hat{\theta} + v_\phi\hat{\phi}$ $\vec{a} = a_r\hat{r} + a_\theta\hat{\theta} + a_\phi\hat{\phi}$
	$v_x(t) = \dot{x}(t)$	$v_r = \dot{r}$	$v_r = \dot{r}$
	$v_y(t) = \dot{y}(t)$	$v_\theta = r\dot{\theta}$	$v_\theta = r\dot{\theta}$
	$v_z(t) = \dot{z}(t)$	$v_z = \dot{z}$	$v_\phi = r \sin \theta \dot{\phi}$
	$a_x(t) = \ddot{x}(t)$	$a_r = \ddot{r} - r\dot{\theta}^2$	$a_r = \ddot{r} - r\dot{\theta}^2 - r \sin^2 \theta \dot{\phi}^2$
	$a_y(t) = \ddot{y}(t)$	$a_\theta = 2\dot{r}\dot{\theta} + r\ddot{\theta}$	$a_\theta = 2\dot{r}\dot{\theta} + r\ddot{\theta} - r \sin \theta \cos \theta \dot{\phi}^2$
	$a_z(t) = \ddot{z}(t)$	$a_z = \ddot{z}$	$a_\phi = 2\dot{r} \sin \theta \dot{\phi} + 2r \cos \theta \dot{\theta} \dot{\phi} + r \sin \theta \ddot{\phi}$

Definitions of Basis Vectors in Different Coordinate Systems

$$\hat{e}_1 = \frac{\nabla u_1}{|\nabla u_1|} = h_1 \nabla u_1$$

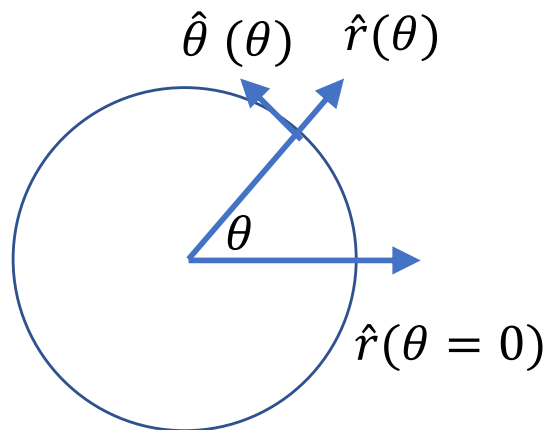
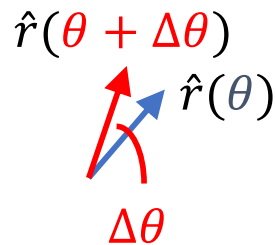
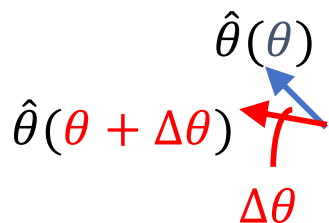
<p>Cartesian coordinate system</p>	$\hat{x} = \frac{\nabla x}{ \nabla x } = \nabla x$	$\hat{y} = \frac{\nabla y}{ \nabla y } = \nabla y$	$\hat{z} = \frac{\nabla z}{ \nabla z } = \nabla z$
<p>Cylindrical coordinate system</p>	$\hat{r} = \frac{\nabla r}{ \nabla r } = \nabla r$	$\hat{\theta} = \frac{\nabla \theta}{ \nabla \theta } = r \nabla \theta$	$\hat{z} = \frac{\nabla z}{ \nabla z } = \nabla z$
<p>Spherical coordinate system</p>	$\hat{r} = \frac{\nabla r}{ \nabla r } = \nabla r$	$\hat{\theta} = \frac{\nabla \theta}{ \nabla \theta } = r \nabla \theta$	$\hat{\phi} = \frac{\nabla \phi}{ \nabla \phi }$ $= r \sin \theta \nabla \phi$

用“幾何”的方式求 單位向量的微分結果

cylindrical coordinate system

$$\frac{d\hat{r}}{d\theta} = \hat{\theta}$$

$$\frac{d\hat{\theta}}{d\theta} = -\hat{r}$$



spherical coordinate system

$$\frac{\partial \hat{r}}{\partial \theta} = \hat{\theta}$$

$$\frac{\partial \hat{\theta}}{\partial \theta} = -\hat{r}$$

$$\frac{\partial \hat{r}}{\partial \phi} = \sin \theta \hat{\phi}$$

$$\frac{\partial \hat{\theta}}{\partial \phi} = \cos \theta \hat{\phi}$$

$$\frac{d\hat{\phi}}{d\phi} = -\sin \theta \hat{r} - \cos \theta \hat{\theta}$$

用“代數”的方式求 單位向量的微分結果

cylindrical coordinate system

$$\hat{r} = \cos \theta \hat{x} + \sin \theta \hat{y}$$

$$\hat{\theta} = -\sin \theta \hat{x} + \cos \theta \hat{y}$$

$$\frac{d\hat{r}}{d\theta} = -\sin \theta \hat{x} + \cos \theta \hat{y} = \hat{\theta}$$

$$\frac{d\hat{\theta}}{d\theta} = -\cos \theta \hat{x} - \sin \theta \hat{y} = -\hat{r}$$

spherical coordinate system

$$\hat{r} = \cos \theta \hat{z} + \sin \theta (\cos \phi \hat{x} + \sin \phi \hat{y})$$

$$\hat{\theta} = -\sin \theta \hat{z} + \cos \theta (\cos \phi \hat{x} + \sin \phi \hat{y})$$

$$\hat{\phi} = -\sin \phi \hat{x} + \cos \phi \hat{y}$$

$$\frac{\partial \hat{r}}{\partial \theta} = -\sin \theta \hat{z} + \cos \theta (\cos \phi \hat{x} + \sin \phi \hat{y}) = \hat{\theta}$$

$$\frac{\partial \hat{r}}{\partial \phi} = \sin \theta (-\sin \phi \hat{x} + \cos \phi \hat{y}) = \sin \theta \hat{\phi}$$

$$\frac{\partial \hat{\theta}}{\partial \theta} = -\cos \theta \hat{z} - \sin \theta (\cos \phi \hat{x} + \sin \phi \hat{y}) = -\hat{r}$$

$$\frac{\partial \hat{\theta}}{\partial \phi} = \cos \theta (-\sin \phi \hat{x} + \cos \phi \hat{y}) = \cos \theta \hat{\phi}$$

$$\frac{d\hat{\phi}}{d\phi} = -\cos \phi \hat{x} - \sin \phi \hat{y}$$

$$-\sin \theta \hat{r} - \cos \theta \hat{\theta}$$

$$= -\sin \theta [\cos \theta \hat{z} + \sin \theta (\cos \phi \hat{x} + \sin \phi \hat{y})]$$

$$- \cos \theta [-\sin \theta \hat{z} + \cos \theta (\cos \phi \hat{x} + \sin \phi \hat{y})]$$

$$= -\cos \phi \hat{x} - \sin \phi \hat{y}$$

$$\frac{d\hat{\phi}}{d\phi} = -\sin \theta \hat{r} - \cos \theta \hat{\theta}$$

Vector Analysis 直角坐標

$$d\mathbf{r} = \hat{x}dx + \hat{y}dy + \hat{z}dz$$

Volume integration of a scalar field, such as the density field $\rho(\mathbf{r}) = \rho(x, y, z)$

$$\iiint [\rho(\mathbf{r})] d\mathbf{r} = \iiint [\rho(x, y, z)] d^3x = \iiint [\rho(x, y, z)] dx dy dz \quad (1)$$

Vector differential operator

$$\nabla = \frac{\partial}{\partial \mathbf{r}} = \hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z}$$

Given a scalar field $f(\mathbf{r}) = f(\vec{r}) = f(x, y, z)$, then the gradient of $f(\mathbf{r})$ is

$$\text{grad } f(\mathbf{r}) = \nabla f(\mathbf{r}) = \hat{x} \frac{\partial f(x, y, z)}{\partial x} + \hat{y} \frac{\partial f(x, y, z)}{\partial y} + \hat{z} \frac{\partial f(x, y, z)}{\partial z} \quad (2)$$

Given a vector field $\mathbf{A}(\mathbf{r}) = \vec{A}(\vec{r}) = \hat{x}A_x(x, y, z) + \hat{y}A_y(x, y, z) + \hat{z}A_z(x, y, z)$, then the divergence of $\mathbf{A}(\mathbf{r})$ is

$$\text{div } \mathbf{A}(\mathbf{r}) = \nabla \cdot \mathbf{A}(\mathbf{r}) = \frac{\partial A_x(x, y, z)}{\partial x} + \frac{\partial A_y(x, y, z)}{\partial y} + \frac{\partial A_z(x, y, z)}{\partial z} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} \quad (3)$$

and the curl of $\mathbf{A}(\mathbf{r})$ is

$$\text{curl } \mathbf{A}(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r}) = \hat{x} \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) + \hat{y} \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) + \hat{z} \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \quad (4)$$

Vector Analysis 柱面坐標

$$d\mathbf{r} = \hat{r}dr + \hat{\theta}rd\theta + \hat{z}dz$$

Volume integration of a scalar field, such as the density field $\rho(\mathbf{r}) = \rho(r, \theta, z)$

$$\iiint [\rho(\mathbf{r})]d\mathbf{r} = \iiint [\rho(r, \theta, z)] d^3x = \iiint [\rho(r, \theta, z)] dr r d\theta dz \quad (5)$$

Vector differential operator

$$\nabla = \frac{\partial}{\partial \mathbf{r}} = \hat{r}\frac{\partial}{\partial r} + \hat{\theta}\frac{1}{r}\frac{\partial}{\partial \theta} + \hat{z}\frac{\partial}{\partial z}$$

Given a scalar field $f(\mathbf{r}) = f(\vec{r}) = f(r, \theta, z)$, then the gradient of $f(\mathbf{r})$ is

$$\text{grad } f(\mathbf{r}) = \nabla f(\mathbf{r}) = \hat{r}\frac{\partial f(r, \theta, z)}{\partial r} + \hat{\theta}\frac{1}{r}\frac{\partial f(r, \theta, z)}{\partial \theta} + \hat{z}\frac{\partial f(r, \theta, z)}{\partial z} \quad (6)$$

Given a vector field $\mathbf{A}(\mathbf{r}) = \vec{A}(\vec{r}) = \hat{r}(\theta) A_r(r, \theta, z) + \hat{\theta}(\theta) A_\theta(r, \theta, z) + \hat{z} A_z(r, \theta, z)$, then the divergence of $\mathbf{A}(\mathbf{r})$ is

$$\text{div } \mathbf{A}(\mathbf{r}) = \nabla \cdot \mathbf{A}(\mathbf{r}) = \frac{\partial A_r}{\partial r} + \frac{A_r}{r} + \frac{1}{r}\frac{\partial A_\theta}{\partial \theta} + \frac{\partial A_z}{\partial z} \quad (7)$$

and the curl of $\mathbf{A}(\mathbf{r})$ is

$$\text{curl } \mathbf{A}(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r}) = \hat{r}\left(\frac{1}{r}\frac{\partial A_z}{\partial \theta} - \frac{\partial A_\theta}{\partial z}\right) + \hat{\theta}\left(\frac{\partial A_r}{\partial z} - \frac{\partial A_z}{\partial r}\right) + \hat{z}\left(\frac{\partial A_\theta}{\partial r} - \frac{1}{r}\frac{\partial A_r}{\partial \theta} + \frac{A_\theta}{r}\right) \quad (8)$$

Vector Analysis 球面坐標

$$d\mathbf{r} = \hat{r} dr + \hat{\theta} r d\theta + \hat{\phi} r \sin \theta d\phi$$

Volume integration of a scalar field, such as the density field $\rho(\mathbf{r}) = \rho(r, \theta, \phi)$

$$\iiint [\rho(\mathbf{r})] d\mathbf{r} = \iiint [\rho(r, \theta, \phi)] d^3x = \iiint [\rho(r, \theta, \phi)] dr r d\theta r \sin \theta d\phi = \iiint [\rho(r, \theta, \phi)] r^2 \sin \theta dr d\theta d\phi \quad (9)$$

Vector differential operator

$$\nabla = \frac{\partial}{\partial \mathbf{r}} = \hat{r} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi}$$

Given a scalar field $f(\mathbf{r}) = f(\vec{r}) = f(r, \theta, \phi)$, then the gradient of $f(\mathbf{r})$ is

$$\text{grad } f(\mathbf{r}) = \nabla f(\mathbf{r}) = \hat{r} \frac{\partial f(r, \theta, \phi)}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial f(r, \theta, \phi)}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial f(r, \theta, \phi)}{\partial \phi} \quad (10)$$

Given a vector field $\mathbf{A}(\mathbf{r}) = \vec{A}(\vec{r}) = \hat{r}(\theta, \phi) A_r(r, \theta, \phi) + \hat{\theta}(\theta, \phi) A_\theta(r, \theta, \phi) + \hat{\phi}(\phi) A_\phi(r, \theta, \phi)$, then the divergence of $\mathbf{A}(\mathbf{r})$ is

$$\text{div } \mathbf{A}(\mathbf{r}) = \nabla \cdot \mathbf{A}(\mathbf{r}) = \frac{\partial A_r}{\partial r} + \frac{2A_r}{r} + \frac{\cos \theta A_\theta}{r \sin \theta} + \frac{1}{r} \frac{\partial A_\theta}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi} \quad (11)$$

and the curl of $\mathbf{A}(\mathbf{r})$ is

$$\text{curl } \mathbf{A}(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r})$$

$$= \hat{r} \left(\frac{1}{r} \frac{\partial A_\phi}{\partial \theta} - \frac{1}{r \sin \theta} \frac{\partial A_\theta}{\partial \phi} + \frac{\cos \theta A_\phi}{r \sin \theta} \right) + \hat{\theta} \left(\frac{1}{r \sin \theta} \frac{\partial A_r}{\partial \phi} - \frac{\partial A_\phi}{\partial r} - \frac{A_\phi}{r} \right) + \hat{\phi} \left(\frac{\partial A_\theta}{\partial r} - \frac{1}{r} \frac{\partial A_r}{\partial \theta} + \frac{A_\theta}{r} \right) \quad (12)$$

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Gauss' theorem & Stokes' theorem

- the divergence theorem, or Gauss' theorem:

微觀的物理量

$$\iiint_{Vol.(S)} \nabla \cdot \vec{E} \, d^3x = \oiint_{S(Vol.)} \vec{E} \cdot d\vec{a}$$

宏觀的物理量

- Stokes' theorem:

$$\iint_{S(C)} (\nabla \times \vec{E}) \cdot d\vec{a} = \oint_{C(S)} \vec{E} \cdot d\vec{l}$$

$$\iint_{S(C)} (\nabla \times \vec{V}) \cdot d\vec{a} = \oint_{C(S)} \vec{V} \cdot d\vec{l}$$

If \vec{V} is the flow field(流場), then $\nabla \times \vec{V}$ is called the vorticity (渦度), and $\oint_{C(S)} \vec{V} \cdot d\vec{l}$ is called the circulation along the closed loop $C(S)$.

以 Stokes' theorem

為範例證明之

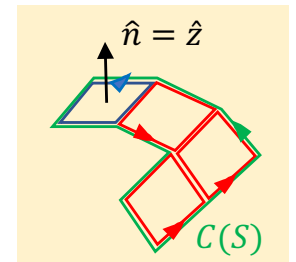
$$\iint_{S(C)} (\nabla \times \vec{E}) \cdot d\vec{a} = \oint_{C(S)} \vec{E} \cdot d\vec{l}$$

將 $S(C)$ 切成很多小塊的方型切面。每塊面的法線方向不一定相同。但是只要每一塊面都滿足上式，則加起來，相鄰兩面右側積分互相抵消，最後只剩最外圍的積分 $C(S)$ 。現在取 local coordinate。考慮其中一個小塊方型切面 S_0 ，法線方向為 \hat{z} 。整塊面位在 $z = 0$ 處。方型切面兩邊分別沿 \hat{x} and \hat{y} 方向。若方型切面左下角位在 (x, y) ，右上角位在 $(x + \Delta x, y + \Delta y)$ 。再令 $(x_0, y_0) = (x + \Delta x/2, y + \Delta y/2)$ 。因為

$$\nabla \times \vec{E} = \hat{x} \left(\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \right) + \hat{y} \left(\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \right) + \hat{z} \left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right)$$

在 $(x_0, y_0, z = 0)$ 處

$$\left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right)_{\substack{x=x_0 \\ y=y_0 \\ z=0}} = \lim_{\substack{\Delta x \rightarrow 0 \\ \Delta y \rightarrow 0}} \left[\frac{E_y(x + \Delta x, y_0) - E_y(x, y_0)}{\Delta x} - \frac{E_x(x_0, y + \Delta y) - E_x(x_0, y)}{\Delta y} \right]_{z=0}$$



$$\begin{aligned} \iint_{S_0} (\nabla \times \vec{E}) \cdot d\vec{a} &= [(\nabla \times \vec{E}) \cdot \hat{z}(\Delta x \Delta y)]_{\substack{x=x_0 \\ y=y_0 \\ z=0}} = \left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right)_{\substack{x=x_0 \\ y=y_0 \\ z=0}} (\Delta x \Delta y) \\ &= \lim_{\substack{\Delta x \rightarrow 0 \\ \Delta y \rightarrow 0}} \left[\frac{E_y(x + \Delta x, y_0) - E_y(x, y_0)}{\Delta x} - \frac{E_x(x_0, y + \Delta y) - E_x(x_0, y)}{\Delta y} \right]_{z=0} (\Delta x \Delta y) \\ &= \lim_{\substack{\Delta x \rightarrow 0 \\ \Delta y \rightarrow 0 \\ z=0}} \{ [E_y(x + \Delta x, y_0) - E_y(x, y_0)] \Delta y - [E_x(x_0, y + \Delta y) - E_x(x_0, y)] \Delta x \} \\ &= \oint_{C(S_0)} \vec{E} \cdot d\vec{l} \end{aligned}$$

