

Appendix B Gauss Theorem
(or Gauss Divergence Theorem)
證明與應用

Show that

$$\oiint_{S(Vol.)} \vec{E} \cdot d\vec{a} = \iiint_{Vol.} \frac{\rho_c}{\epsilon_0} d^3x$$

yields

$$\nabla \cdot \vec{E} = \frac{\rho_c}{\epsilon_0}$$

and vice versa (the other way around).

Let us consider an infinitesimal volume $Vol. = \Delta x \Delta y \Delta z$ centered at (x_0, y_0, z_0) . The charge density inside this volume can be considered as a uniform function, thus we have

$$\iiint_{Vol.} \frac{\rho_c}{\epsilon_0} d^3x \approx \frac{\rho_c(x_0, y_0, z_0)}{\epsilon_0} \Delta x \Delta y \Delta z \quad (\text{B.1})$$

and the surface around the infinitesimal volume $Vol. = \Delta x \Delta y \Delta z$ should locate at $x = x_0 \pm \Delta x/2$, $y = y_0 \pm \Delta y/2$, and $z = z_0 \pm \Delta z/2$. Thus

$$\begin{aligned}
& \oiint_{S(\text{Vol.})} \vec{E} \cdot d\vec{a} \\
&= E_x \left(x_0 + \frac{\Delta x}{2}, y_0, z_0 \right) \Delta y \Delta z \\
&\quad - E_x \left(x_0 - \frac{\Delta x}{2}, y_0, z_0 \right) \Delta y \Delta z \\
&\quad + E_y \left(x_0, y_0 + \frac{\Delta y}{2}, z_0 \right) \Delta x \Delta z \\
&\quad - E_y \left(x_0, y_0 - \frac{\Delta y}{2}, z_0 \right) \Delta x \Delta z \\
&\quad + E_z \left(x_0, y_0, z_0 + \frac{\Delta z}{2} \right) \Delta x \Delta y \\
&\quad - E_z \left(x_0, y_0, z_0 - \frac{\Delta z}{2} \right) \Delta x \Delta y
\end{aligned} \tag{B.2}$$

Substituting Equation (B.1) and (B.2) into the following equation

$$\oiint_{S(Vol.)} \vec{E} \cdot d\vec{a} = \iiint_{Vol.} \frac{\rho_c}{\epsilon_0} d^3x$$

It yields

$$\begin{aligned} & E_x \left(x_0 + \frac{\Delta x}{2}, y_0, z_0 \right) \Delta y \Delta z - E_x \left(x_0 - \frac{\Delta x}{2}, y_0, z_0 \right) \Delta y \Delta z \\ & + E_y \left(x_0, y_0 + \frac{\Delta y}{2}, z_0 \right) \Delta x \Delta z - E_y \left(x_0, y_0 - \frac{\Delta y}{2}, z_0 \right) \Delta x \Delta z \\ & + E_z \left(x_0, y_0, z_0 + \frac{\Delta z}{2} \right) \Delta x \Delta y - E_z \left(x_0, y_0, z_0 - \frac{\Delta z}{2} \right) \Delta x \Delta y \\ & = \frac{\rho_c(x_0, y_0, z_0)}{\epsilon_0} \Delta x \Delta y \Delta z \end{aligned}$$

Multiplying the above equation by $1/\Delta x\Delta y\Delta z$, it yields

$$\begin{aligned} & \frac{E_x\left(x_0 + \frac{\Delta x}{2}, y_0, z_0\right) - E_x\left(x_0 - \frac{\Delta x}{2}, y_0, z_0\right)}{\Delta x} \\ & + \frac{E_y\left(x_0, y_0 + \frac{\Delta y}{2}, z_0\right) - E_y\left(x_0, y_0 - \frac{\Delta y}{2}, z_0\right)}{\Delta y} \\ & + \frac{E_z\left(x_0, y_0, z_0 + \frac{\Delta z}{2}\right) - E_z\left(x_0, y_0, z_0 - \frac{\Delta z}{2}\right)}{\Delta z} \\ & = \frac{\rho_c(x_0, y_0, z_0)}{\epsilon_0} \end{aligned}$$

For an infinitesimal vol. $\Delta x \Delta y \Delta z$, it yields

$$\left[\nabla \cdot \vec{E} = \frac{\rho_c}{\epsilon_0} \right]_{(x_0, y_0, z_0)} \quad Q.E.D.$$

Likewise, for

$$\nabla \cdot \vec{E} = \frac{\rho_c}{\epsilon_0}$$

it yields

$$\oiint_{S(\text{Vol.})} \vec{E} \cdot d\vec{a} = \iiint_{\text{Vol.}} \frac{\rho_c}{\epsilon_0} d^3x$$

in an infinitesimal vol. $\Delta x \Delta y \Delta z$.

把所有這樣的小體積都加起來，就是總體積，而相鄰的「面積分」會互相抵銷，只剩下最外圍的「封閉曲面積分」。故得證

$\nabla \cdot \vec{E} = \rho_c/\epsilon_0$ ，可推得以下積分形式

$$\oiint_{S(\text{Vol.})} \vec{E} \cdot d\vec{a} = \iiint_{\text{Vol.}} \frac{\rho_c}{\epsilon_0} d^3x$$

Likewise, $\nabla \cdot \vec{B} = 0$ yields

$$\oiint_{S(\text{Vol.})} \vec{B} \cdot d\vec{a} = 0$$

由此結果可以證明磁場線密集處磁場比較強。(請證明)

Likewise, the continuity equation

$$\left[\frac{\partial}{\partial t} + \vec{V}_\alpha \cdot \nabla \right] n_\alpha = -n_\alpha \nabla \cdot \vec{V}_\alpha$$

yields, for incompressible fluid, 流線密集處, 流速比較大,
where incompressible fluid is defined by

$$\left[\frac{\partial}{\partial t} + \vec{V}_\alpha \cdot \nabla \right] n_\alpha = \frac{dn_\alpha}{dt} \Big|_{\text{along the trajectory of a fluid element}} = 0$$